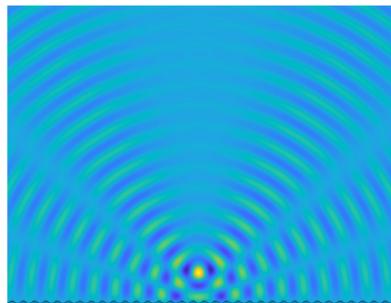
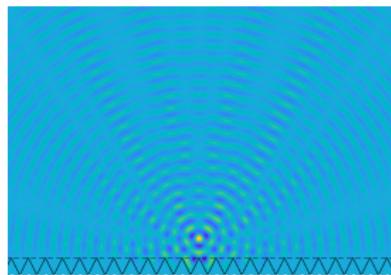


# Integral equations and boundary element methods for rough surface scattering (RSS)



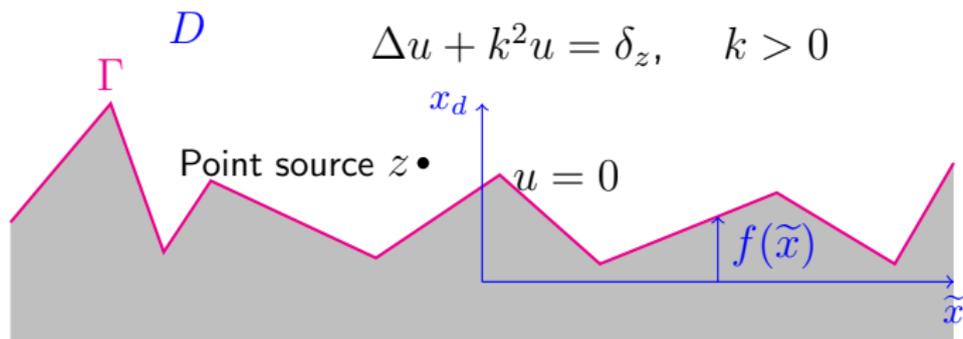
Simon Chandler-Wilde

Department of Mathematics  
and Statistics  
University of Reading, UK

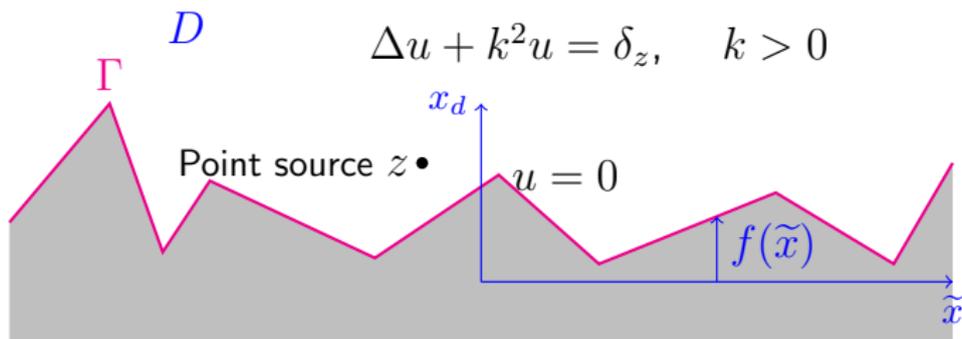


With: Martin Averseng & Euan Spence (Bath, UK)

INI Computational Methods for Multiple Scattering Workshop, April 2023

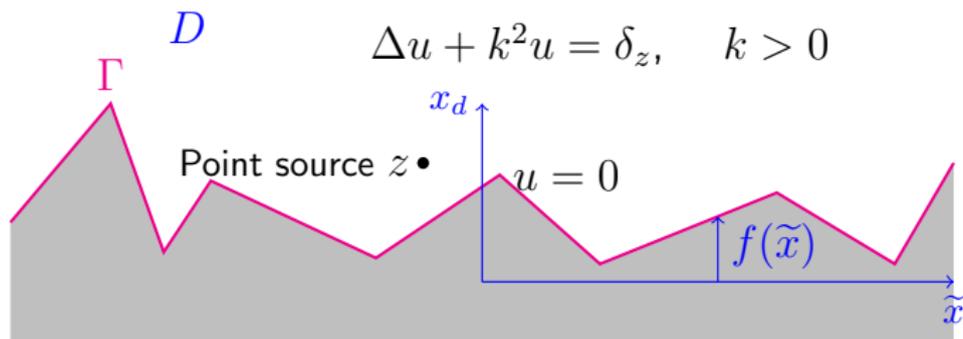


Many interesting **computational and numerical analysis** challenges!



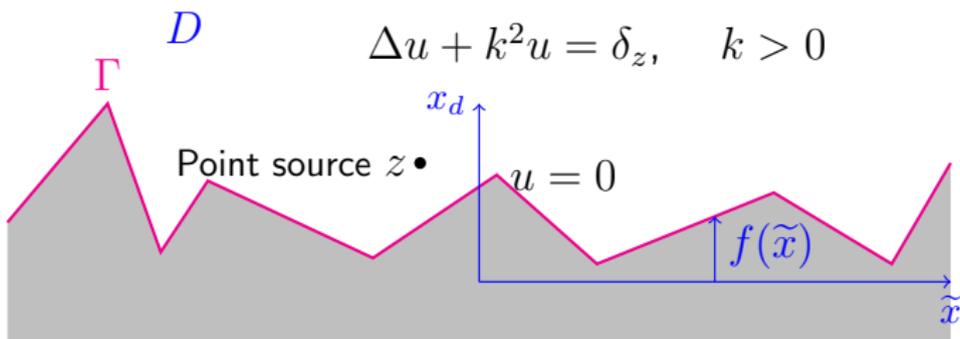
Many interesting **computational and numerical analysis challenges!**

- Formulation, including radiation condition, and well-posedness, clear in the above case (C-W, Heinmeyer, Potthast, 2006, C-W, Elschner, 2010), but ...



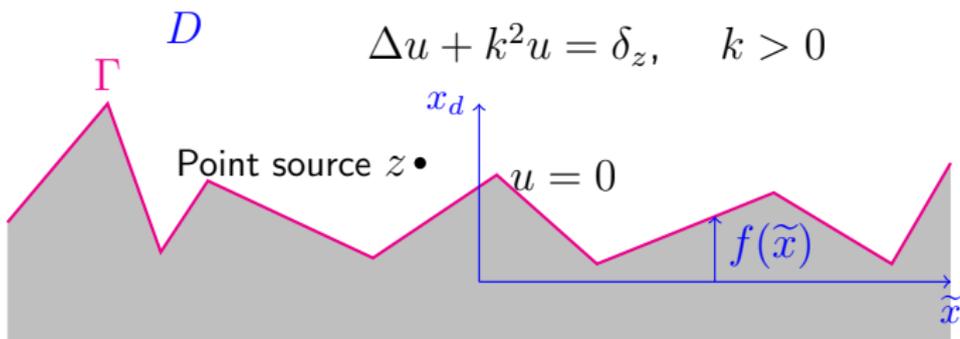
Many interesting **computational and numerical analysis challenges!**

- Formulation, including radiation condition, and well-posedness, clear in the above case (C-W, Heinmeyer, Potthast, 2006, C-W, Elschner, 2010), but ...
- Non-uniqueness - solutions to homogeneous problem localised near  $\Gamma$  for Neumann b.c. or if  $\Gamma$  not a graph (Gotlib, 2000)



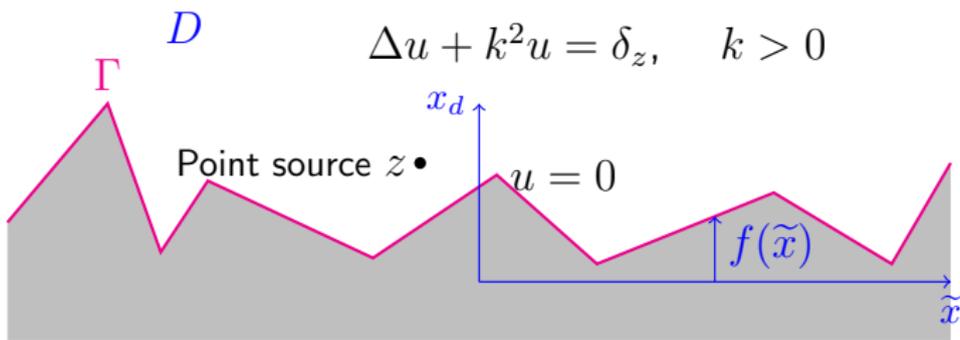
Many interesting **computational and numerical analysis challenges!**

- Formulation, including radiation condition, and well-posedness, clear in the above case (C-W, Heinmeyer, Potthast, 2006, C-W, Elschner, 2010), but ...
- Non-uniqueness - solutions to homogeneous problem localised near  $\Gamma$  for Neumann b.c. or if  $\Gamma$  not a graph (Gotlib, 2000)
- Unclear whether plane wave incidence makes sense in general in 3D (see Rathsfeld 2022.)



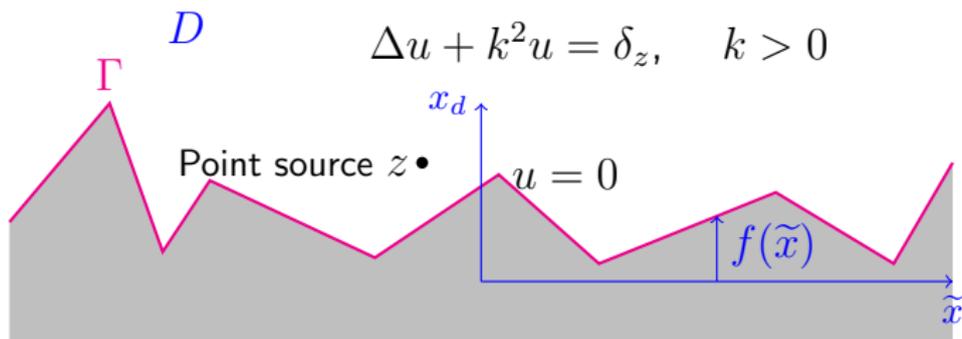
Many interesting **computational and numerical analysis challenges!**

- Formulation, including radiation condition, and well-posedness, clear in the above case (C-W, Heinmeyer, Potthast, 2006, C-W, Elschner, 2010), but ...
- Non-uniqueness - solutions to homogeneous problem localised near  $\Gamma$  for Neumann b.c. or if  $\Gamma$  not a graph (Gotlib, 2000)
- Unclear whether plane wave incidence makes sense in general in 3D (see Rathsfeld 2022.)
- Usual boundary integral equations (BIE) methods for bounded obstacles very popular, but:
  - i) need to discretize large section of  $\Gamma$  of diameter  $2a$  for accuracy;
  - ii) condition numbers for standard methods grow at least like  $(ka)^{1/2}$



Many interesting **computational and numerical analysis challenges!**

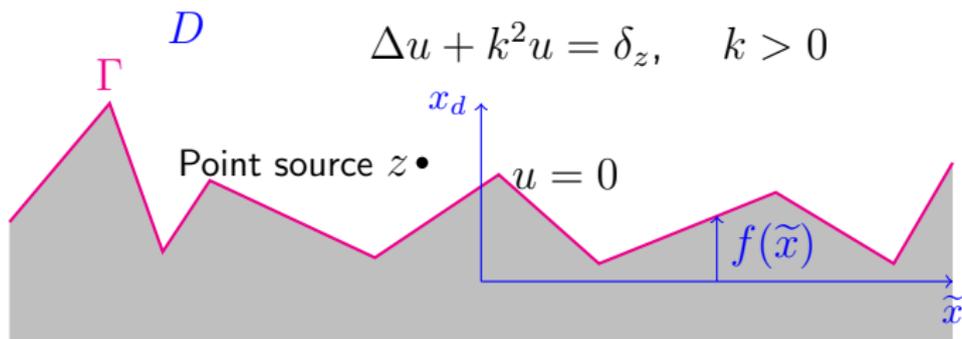
- Formulation, including radiation condition, and well-posedness, clear in the above case (C-W, Heinmeyer, Potthast, 2006, C-W, Elschner, 2010), but ...
- Non-uniqueness - solutions to homogeneous problem localised near  $\Gamma$  for Neumann b.c. or if  $\Gamma$  not a graph (Gotlib, 2000)
- Unclear whether plane wave incidence makes sense in general in 3D (see Rathsfeld 2022.)
- Usual boundary integral equations (BIE) methods for bounded obstacles very popular, but:
  - i) need to discretize large section of  $\Gamma$  of diameter  $2a$  for accuracy;
  - ii) condition numbers for standard methods grow at least like  $(ka)^{1/2}$
- Numerical analysis challenges: stability and convergence of truncation of unbounded surface? Analysis of boundary element methods (BEM) when surface is unbounded, and of convergence of iterative solvers (GMRES)?



Many interesting **computational and numerical analysis challenges!**

- Usual boundary integral equations (BIE) methods for bounded obstacles very popular, but:
  - i) need to discretize large section of  $\Gamma$  of diameter  $2a$  for accuracy;
  - ii) condition numbers for standard methods grow at least like  $(ka)^{1/2}$
- Numerical analysis challenges: stability and convergence of truncation of unbounded surface; analysis of boundary element methods (BEM) when surface is unbounded, and of convergence of iterative solvers (GMRES).

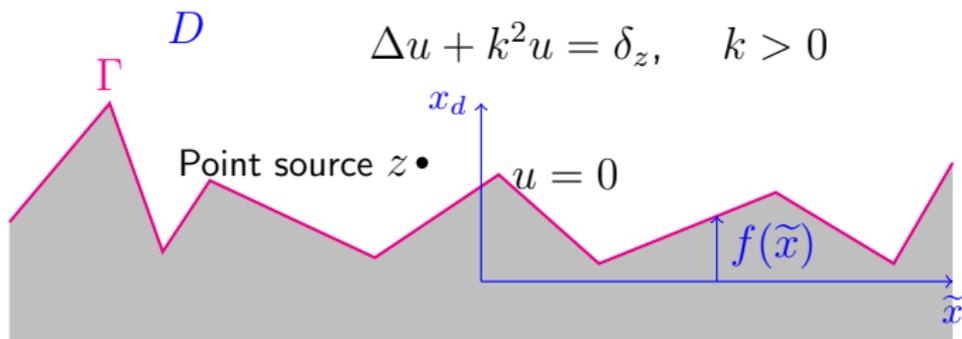
In this talk we will:



Many interesting **computational and numerical analysis challenges!**

- Usual boundary integral equations (BIE) methods for bounded obstacles very popular, but:
  - i) need to discretize large section of  $\Gamma$  of diameter  $2a$  for accuracy;
  - ii) condition numbers for standard methods grow at least like  $(ka)^{1/2}$
- Numerical analysis challenges: stability and convergence of truncation of unbounded surface; analysis of boundary element methods (BEM) when surface is unbounded, and of convergence of iterative solvers (GMRES).

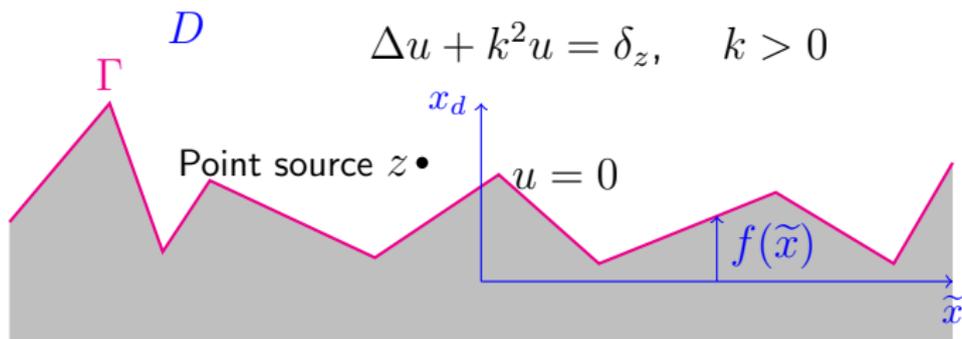
In this talk we will: **propose a new 2nd kind BIE** for the above problem with operator that is **uniformly bounded in  $a$  and coercive with coercivity constant dependent only on the maximum surface slope;**



Many interesting **computational and numerical analysis challenges!**

- Usual boundary integral equations (BIE) methods for bounded obstacles very popular, but:
  - i) need to discretize large section of  $\Gamma$  of diameter  $2a$  for accuracy;
  - ii) condition numbers for standard methods grow at least like  $(ka)^{1/2}$
- Numerical analysis challenges: stability and convergence of truncation of unbounded surface; analysis of boundary element methods (BEM) when surface is unbounded, and of convergence of iterative solvers (GMRES).

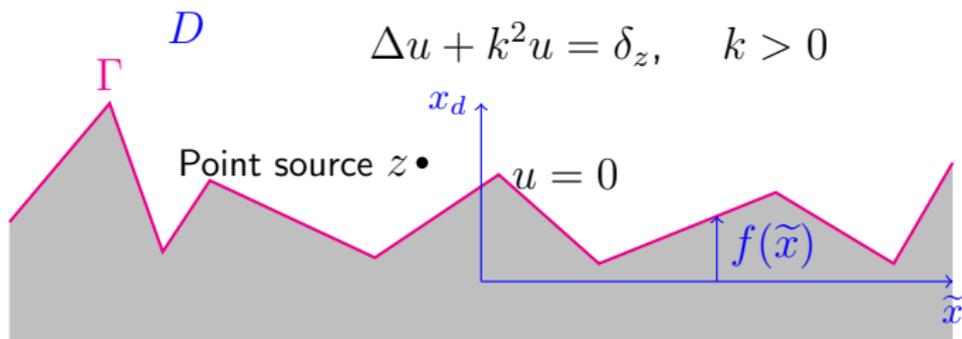
In this talk we will: **propose a new 2nd kind BIE** for the above problem with operator that is **uniformly bounded in  $a$  and coercive with coercivity constant dependent only on the maximum surface slope**; **prove convergence of combined Galerkin BEM/ surface truncation**;



Many interesting **computational and numerical analysis challenges!**

- Usual boundary integral equations (BIE) methods for bounded obstacles very popular, but:
  - i) need to discretize large section of  $\Gamma$  of diameter  $2a$  for accuracy;
  - ii) condition numbers for standard methods grow at least like  $(ka)^{1/2}$
- Numerical analysis challenges: stability and convergence of truncation of unbounded surface; analysis of boundary element methods (BEM) when surface is unbounded, and of convergence of iterative solvers (GMRES).

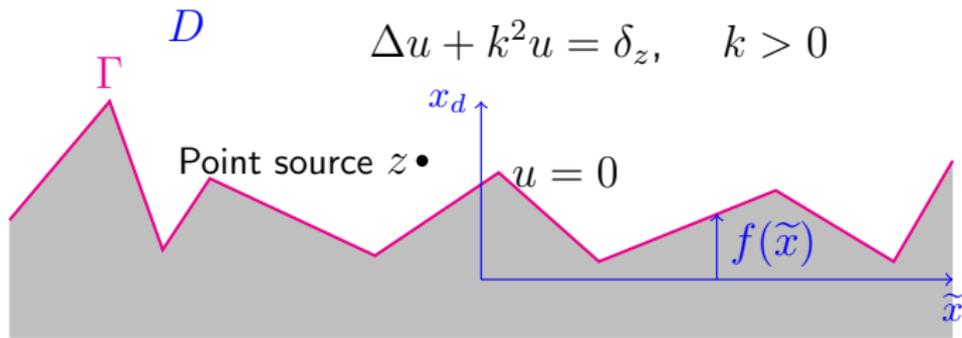
In this talk we will: **propose a new 2nd kind BIE** for the above problem with operator that is **uniformly bounded in  $a$  and coercive with coercivity constant dependent only on the maximum surface slope**; prove **convergence of combined Galerkin BEM/ surface truncation**; prove that a **fixed number of GMRES iterations is sufficient, uniformly in the BEM step size ( $h$ ) and the size of the truncated surface discretized ( $a$ )**.



Many interesting **computational and numerical analysis challenges!**

- Usual boundary integral equations (BIE) methods for bounded obstacles very popular, but:
  - i) need to discretize large section of  $\Gamma$  of diameter  $2a$  for accuracy;
  - ii) condition numbers for standard methods grow at least like  $(ka)^{1/2}$
- Numerical analysis challenges: stability and convergence of truncation of unbounded surface; analysis of boundary element methods (BEM) when surface is unbounded, and of convergence of iterative solvers (GMRES).

In this talk we will: **propose a new 2nd kind BIE** for the above problem with operator that is **uniformly bounded in  $a$  and coercive with coercivity constant dependent only on the maximum surface slope**; prove **convergence of combined Galerkin BEM/ surface truncation**; prove that a **fixed number of GMRES iterations is sufficient, uniformly in the BEM step size ( $h$ ) and the size of the truncated surface discretized ( $a$ )**. On the way we will recall: existing analysis tools for Galerkin BEM/GMRES;



Many interesting **computational and numerical analysis challenges!**

- Usual boundary integral equations (BIE) methods for bounded obstacles very popular, but:
  - i) need to discretize large section of  $\Gamma$  of diameter  $2a$  for accuracy;
  - ii) condition numbers for standard methods grow at least like  $(ka)^{1/2}$
- Numerical analysis challenges: stability and convergence of truncation of unbounded surface; analysis of boundary element methods (BEM) when surface is unbounded, and of convergence of iterative solvers (GMRES).

In this talk we will: **propose a new 2nd kind BIE** for the above problem with operator that is **uniformly bounded in  $a$  and coercive with coercivity constant dependent only on the maximum surface slope**; prove **convergence of combined Galerkin BEM/surface truncation**; prove that a **fixed number of GMRES iterations is sufficient, uniformly in the BEM step size ( $h$ ) and the size of the truncated surface discretized ( $a$ )**. On the way we will recall: existing analysis tools for Galerkin BEM/GMRES; recent related results for 2nd kind BIEs for (single and multiple) bounded scatterers.

# Tools for convergence of Galerkin methods and GMRES

Suppose that  $H$  is a complex Hilbert space with norm  $\|u\| = \sqrt{(u, u)}$ , e.g.

$$H = L^2(\Gamma), \quad (u, v) = \int_{\Gamma} u \bar{v} \, ds, \quad \|u\|^2 = \int_{\Gamma} |u|^2 \, ds.$$

# Tools for convergence of Galerkin methods and GMRES

Suppose that  $H$  is a complex Hilbert space with norm  $\|u\| = \sqrt{(u, u)}$ , e.g.

$$H = L^2(\Gamma), \quad (u, v) = \int_{\Gamma} u \bar{v} \, ds, \quad \|u\|^2 = \int_{\Gamma} |u|^2 \, ds.$$

Suppose that  $A$  is a **bounded linear operator** on  $H$  and that  $A$  is **coercive**, i.e., for some  $\gamma > 0$ ,

$$|(Au, u)| \geq \gamma \|u\|^2, \quad \forall u \in H.$$

# Tools for convergence of Galerkin methods and GMRES

Suppose that  $H$  is a complex Hilbert space with norm  $\|u\| = \sqrt{(u, u)}$ , e.g.

$$H = L^2(\Gamma), \quad (u, v) = \int_{\Gamma} u \bar{v} \, ds, \quad \|u\|^2 = \int_{\Gamma} |u|^2 \, ds.$$

Suppose that  $A$  is a **bounded linear operator** on  $H$  and that  $A$  is **coercive**, i.e., for some  $\gamma > 0$ ,

$$|(Au, u)| \geq \gamma \|u\|^2, \quad \forall u \in H.$$

**Lax-Milgram Lemma.**  $A$  is invertible and  $\|A^{-1}\| \leq \gamma^{-1}$ .

# Tools for convergence of Galerkin methods and GMRES

Suppose that  $H$  is a complex Hilbert space with norm  $\|u\| = \sqrt{(u, u)}$ , e.g.

$$H = L^2(\Gamma), \quad (u, v) = \int_{\Gamma} u \bar{v} \, ds, \quad \|u\|^2 = \int_{\Gamma} |u|^2 \, ds.$$

Suppose that  $A$  is a **bounded linear operator** on  $H$  and that  $A$  is **coercive**, i.e., for some  $\gamma > 0$ ,

$$|(Au, u)| \geq \gamma \|u\|^2, \quad \forall u \in H.$$

**Lax-Milgram Lemma.**  $A$  is invertible and  $\|A^{-1}\| \leq \gamma^{-1}$ .

**Céa's Lemma.** Let  $H_N \subset H$  be a closed subspace. Then,  $\forall g \in H$ ,  $\exists$  a unique *Galerkin approximation*  $u_N \in H_N$  to  $u := A^{-1}g$ , defined by

$$(Au_N, v_N) = (g, v), \quad \forall v_N \in H_N,$$

and

$$\|u - u_N\| \leq \frac{\|A\|}{\gamma} \inf_{v_N \in H_N} \|u - v_N\|.$$

# Tools for convergence of Galerkin methods and GMRES

Suppose that  $H$  is a complex Hilbert space with norm  $\|u\| = \sqrt{(u, u)}$ , e.g.

$$H = L^2(\Gamma), \quad (u, v) = \int_{\Gamma} u \bar{v} \, ds, \quad \|u\|^2 = \int_{\Gamma} |u|^2 \, ds.$$

Suppose that  $A$  is a **bounded linear operator** on  $H$  and that  $A$  is **coercive**, i.e., for some  $\gamma > 0$ ,

$$|(Au, u)| \geq \gamma \|u\|^2, \quad \forall u \in H.$$

**Lax-Milgram Lemma.**  $A$  is invertible and  $\|A^{-1}\| \leq \gamma^{-1}$ .

**Céa's Lemma.** Let  $H_N \subset H$  be a closed subspace. Then,  $\forall g \in H$ ,  $\exists$  a unique *Galerkin approximation*  $u_N \in H_N$  to  $u := A^{-1}g$ , defined by

$$(Au_N, v_N) = (g, v), \quad \forall v_N \in H_N,$$

and

$$\|u - u_N\| \leq \frac{\|A\|}{\gamma} \inf_{v_N \in H_N} \|u - v_N\|. \quad \text{Note } \frac{\|A\|}{\gamma} \geq \text{cond}(A) := \|A\| \|A\|^{-1}.$$

Suppose that  $A$  is a **bounded linear operator** on  $H$  and that  $A$  is **coercive**, i.e., for some  $\gamma > 0$ ,

$$|(Au, u)| \geq \gamma \|u\|^2, \quad \forall u \in H.$$

**Céa's Lemma.** Let  $H_N \subset H$  be a closed subspace. Then,  $\forall g \in H$ ,  $\exists$  a unique *Galerkin approximation*  $u_N \in H_N$  to  $u := A^{-1}g$ , defined by

$$(Au_N, v_N) = (g, v), \quad \forall v_N \in H_N, \quad (*)$$

and

$$\|u - u_N\| \leq \frac{\|A\|}{\gamma} \inf_{v_N \in H_N} \|u - v_N\|.$$

Suppose that  $A$  is a **bounded linear operator** on  $H$  and that  $A$  is **coercive**, i.e., for some  $\gamma > 0$ ,

$$|(Au, u)| \geq \gamma \|u\|^2, \quad \forall u \in H.$$

**Céa's Lemma.** Let  $H_N \subset H$  be a closed subspace. Then,  $\forall g \in H, \exists$  a unique *Galerkin approximation*  $u_N \in H_N$  to  $u := A^{-1}g$ , defined by

$$(Au_N, v_N) = (g, v_N), \quad \forall v_N \in H_N, \quad (*)$$

and

$$\|u - u_N\| \leq \frac{\|A\|}{\gamma} \inf_{v_N \in H_N} \|u - v_N\|.$$

If  $H_N$  has basis  $\{\varphi_1, \dots, \varphi_N\}$ , then  $u_N = \sum_{n=1}^N a_n \varphi_n$  and  $(*)$  is

$$\sum_{n=1}^N (A\varphi_n, \varphi_m) a_n = (g, \varphi_m), \quad m = 1, \dots, N. \quad (X)$$

Suppose that  $A$  is a **bounded linear operator** on  $H$  and that  $A$  is **coercive**, i.e., for some  $\gamma > 0$ ,

$$|(Au, u)| \geq \gamma \|u\|^2, \quad \forall u \in H.$$

**Céa's Lemma.** Let  $H_N \subset H$  be a closed subspace. Then,  $\forall g \in H$ ,  $\exists$  a unique *Galerkin approximation*  $u_N \in H_N$  to  $u := A^{-1}g$ , defined by

$$(Au_N, v_N) = (g, v_N), \quad \forall v_N \in H_N, \quad (*)$$

and

$$\|u - u_N\| \leq \frac{\|A\|}{\gamma} \inf_{v_N \in H_N} \|u - v_N\|.$$

If  $H_N$  has basis  $\{\varphi_1, \dots, \varphi_N\}$ , then  $u_N = \sum_{n=1}^N a_n \varphi_n$  and  $(*)$  is

$$\sum_{n=1}^N (A\varphi_n, \varphi_m) a_n = (g, \varphi_m), \quad m = 1, \dots, N. \quad (X)$$

**Theorem** (*corollary of field of values estimate in Beckermann et al. 2006*). Let  $r_m$  be the residual after  $m$  steps of GMRES applied to  $(X)$ . Then

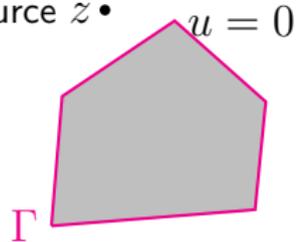
$$\frac{\|r_m\|_2}{\|r_0\|_2} \leq \varepsilon \quad \text{provided} \quad m \geq \frac{3\sqrt{3}}{4} \frac{\|A\|}{\gamma} \text{cond}(M) \log\left(\frac{8}{\varepsilon}\right),$$

where  $M = [(\varphi_n, \varphi_m)]$  is the mass matrix.

# Integral equation methods: bounded obstacle case

Point source  $z \bullet$

$D$

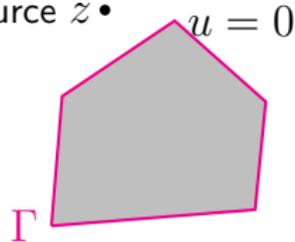


$$\Delta u + k^2 u = \delta_z, \quad k > 0$$

# Integral equation methods: bounded obstacle case

Point source  $z \bullet$

$D$



$$\Delta u + k^2 u = \delta_z, \quad k > 0$$

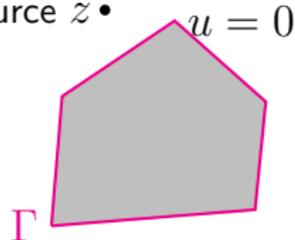
By Green's theorem, where  $\Phi(x, y) := \frac{e^{ik|x-y|}}{4\pi|x-y|}$  and  $u^i(x) := \Phi(x, z)$ ,

$$u(x) = u^i(x) - \int_{\Gamma} \Phi(x, y) \partial_n u(y) \, ds(y), \quad x \in D.$$

# Integral equation methods: bounded obstacle case

Point source  $z \bullet$

$D$



$$\Delta u + k^2 u = \delta_z, \quad k > 0$$

By Green's theorem, where  $\Phi(x, y) := \frac{e^{ik|x-y|}}{4\pi|x-y|}$  and  $u^i(x) := \Phi(x, z)$ ,

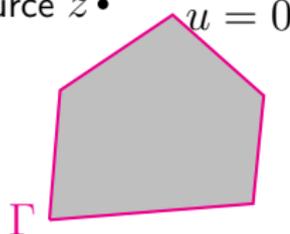
$$u(x) = u^i(x) - \int_{\Gamma} \Phi(x, y) \partial_n u(y) ds(y), \quad x \in D.$$

Taking a linear combination of Dirichlet ( $\gamma$ ) and Neumann ( $\partial_n$ ) traces, we obtain the standard 2nd kind integral equation

# Integral equation methods: bounded obstacle case

Point source  $z \bullet$

$D$



$$\Delta u + k^2 u = \delta_z, \quad k > 0$$

By Green's theorem, where  $\Phi(x, y) := \frac{e^{ik|x-y|}}{4\pi|x-y|}$  and  $u^i(x) := \Phi(x, z)$ ,

$$u(x) = u^i(x) - \int_{\Gamma} \Phi(x, y) \partial_n u(y) ds(y), \quad x \in D.$$

Taking a linear combination of Dirichlet ( $\gamma$ ) and Neumann ( $\partial_n$ ) traces, we obtain the standard 2nd kind integral equation

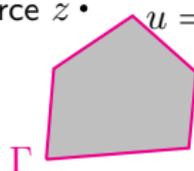
$$A \partial_n u = g := \partial_n u^i - ik \gamma u^i, \quad \text{where} \quad A := \frac{1}{2} I + K' - ik S,$$

$$K' \varphi(x) := \int_{\Gamma} \partial_{n(x)} \Phi(x, y) \varphi(y) ds(y), \quad S \varphi(x) := \int_{\Gamma} \Phi(x, y) \varphi(y) ds(y), \quad x \in \Gamma.$$

# Integral equation methods: bounded obstacle case

Point source  $z \bullet$

$D$



$$\Delta u + k^2 u = \delta_z, \quad k > 0$$

$$u(x) = u^i(x) - \int_{\Gamma} \Phi(x, y) \partial_n u(y) \, ds(y), \quad x \in D.$$

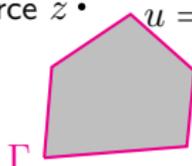
with

$$A \partial_n u = g := \partial_n u^i - ik\gamma u^i, \quad \text{and} \quad A := \frac{1}{2}I + K' - ikS.$$

# Integral equation methods: bounded obstacle case

Point source  $z \bullet$

$D$



$u = 0$

$$\Delta u + k^2 u = \delta_z, \quad k > 0$$

$$u(x) = u^i(x) - \int_{\Gamma} \Phi(x, y) \partial_n u(y) \, ds(y), \quad x \in D.$$

with

$$A \partial_n u = g := \partial_n u^i - ik\gamma u^i, \quad \text{and} \quad A := \frac{1}{2}I + K' - ikS.$$

- $A$  is invertible on  $L^2(\Gamma)$  for general Lipschitz  $\Gamma$  (C-W & Langdon 2007)

# Integral equation methods: bounded obstacle case

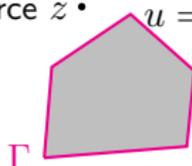
Point source  $z \bullet$

$D$

$\Gamma$

$u = 0$

$\Delta u + k^2 u = \delta_z, \quad k > 0$



$$u(x) = u^i(x) - \int_{\Gamma} \Phi(x, y) \partial_n u(y) \, ds(y), \quad x \in D.$$

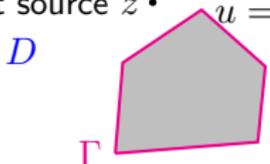
with

$$A \partial_n u = g := \partial_n u^i - ik\gamma u^i, \quad \text{and} \quad A := \frac{1}{2}I + K' - ikS.$$

- $A$  is invertible on  $L^2(\Gamma)$  for general Lipschitz  $\Gamma$  (C-W & Langdon 2007)
- $\|A^{-1}\| = O(1)$  as  $k \rightarrow \infty$  if  $\Gamma$  is star-shaped or smooth and non-trapping (C-W & Monk 2008, Baskin, Spence, Wunsch 2016)
- $\|A\| = O(k^{1/2})$  as  $k \rightarrow \infty$  if  $\Gamma$  is star-shaped and 2D, or smooth and non-trapping (C-W et al. 2009, Baskin, Spence, Wunsch 2016)

# Integral equation methods: bounded obstacle case

Point source  $z \bullet$



The diagram shows a gray-shaded polygonal domain  $D$  with a pink boundary  $\Gamma$ . A point source  $z$  is marked with a black dot inside the domain. The boundary is labeled  $u=0$  at the top vertex.

$$\Delta u + k^2 u = \delta_z, \quad k > 0$$

$$u(x) = u^i(x) - \int_{\Gamma} \Phi(x, y) \partial_n u(y) ds(y), \quad x \in D.$$

with

$$A \partial_n u = g := \partial_n u^i - ik\gamma u^i, \quad \text{and} \quad A := \frac{1}{2}I + K' - ikS.$$

- $\|A^{-1}\| = O(1)$  as  $k \rightarrow \infty$  if  $\Gamma$  is star-shaped or smooth and non-trapping (C-W & Monk 2008, Baskin, Spence, Wunsch 2016)
- $\|A\| = O(k^{1/2})$  as  $k \rightarrow \infty$  if  $\Gamma$  is star-shaped and 2D, or smooth and non-trapping (C-W et al. 2009, Baskin, Spence, Wunsch 2016)
- $A$  is **uniformly-in- $k$  coercive**, i.e., for all  $k_0 > 0$  there exists  $\gamma > 0$  such that

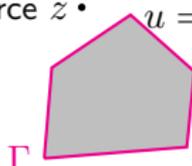
$$|(A\varphi, \varphi)| \geq \gamma \|\varphi\|^2, \quad \varphi \in L^2(\Gamma), \quad k \geq k_0,$$

if  $\Gamma$  is **smooth and uniformly convex** (Spence, Kamotski, Smyshlyaev 2016)

# Integral equation methods: bounded obstacle case

Point source  $z \bullet$

$D$



$u = 0$

$\Gamma$

$\Delta u + k^2 u = \delta_z, \quad k > 0$

$$u(x) = u^i(x) - \int_{\Gamma} \Phi(x, y) \partial_n u(y) ds(y), \quad x \in D.$$

with

$$A \partial_n u = g := \partial_n u^i - ik\gamma u^i, \quad \text{and} \quad A := \frac{1}{2}I + K' - ikS.$$

- $\|A^{-1}\| = O(1)$  as  $k \rightarrow \infty$  if  $\Gamma$  is star-shaped or smooth and non-trapping (C-W & Monk 2008, Baskin, Spence, Wunsch 2016)
- $\|A\| = O(k^{1/2})$  as  $k \rightarrow \infty$  if  $\Gamma$  is star-shaped and 2D, or smooth and non-trapping (C-W et al. 2009, Baskin, Spence, Wunsch 2016)
- $A$  is **uniformly-in- $k$  coercive**, i.e., for all  $k_0 > 0$  there exists  $\gamma > 0$  such that

$$|(A\varphi, \varphi)| \geq \gamma \|\varphi\|^2, \quad \varphi \in L^2(\Gamma), \quad k \geq k_0,$$

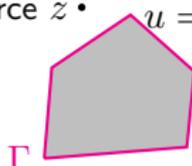
if  $\Gamma$  is **smooth and uniformly convex** (Spence, Kamotski, Smyshlyaev 2016)

- **BUT**  $A$  is not even compactly perturbed coercive for general Lipschitz  $\Gamma$ , or even for general star-shaped polyhedra in 3D (C-W & Spence 2022a)

# Integral equation methods: bounded obstacle case

Point source  $z \bullet$

$D$



$u = 0$

$\Gamma$

$\Delta u + k^2 u = \delta_z, \quad k > 0$

$$u(x) = u^i(x) - \int_{\Gamma} \Phi(x, y) \partial_n u(y) ds(y), \quad x \in D.$$

with

$$A \partial_n u = g := \partial_n u^i - ik\gamma u^i, \quad \text{and} \quad A := \frac{1}{2}I + K' - ikS.$$

- $\|A^{-1}\| = O(1)$  as  $k \rightarrow \infty$  if  $\Gamma$  is star-shaped or smooth and non-trapping (C-W & Monk 2008, Baskin, Spence, Wunsch 2016)
- $\|A\| = O(k^{1/2})$  as  $k \rightarrow \infty$  if  $\Gamma$  is star-shaped and 2D, or smooth and non-trapping (C-W et al. 2009, Baskin, Spence, Wunsch 2016)
- $A$  is **uniformly-in- $k$  coercive**, i.e., for all  $k_0 > 0$  there exists  $\gamma > 0$  such that

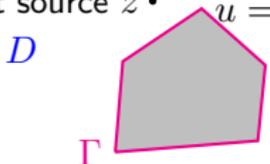
$$|(A\varphi, \varphi)| \geq \gamma \|\varphi\|^2, \quad \varphi \in L^2(\Gamma), \quad k \geq k_0,$$

if  $\Gamma$  is **smooth and uniformly convex** (Spence, Kamotski, Smyshlyaev 2016)

- **BUT**  $A$  is not even compactly perturbed coercive for general Lipschitz  $\Gamma$ , or even for general star-shaped polyhedra in 3D (C-W & Spence 2022a) **AND** there is no numerical method provably convergent for every polyhedron  $\Gamma$  (**open problem**).

## Coercive formulations: bounded obstacle case

Point source  $z \bullet$



The diagram shows a gray shaded pentagonal domain  $D$  with a pink boundary  $\Gamma$ . A point source  $z$  is marked with a black dot inside the domain. The boundary  $\Gamma$  is labeled with  $u = 0$  at the top vertex.

$$\Delta u + k^2 u = \delta_z, \quad k > 0$$

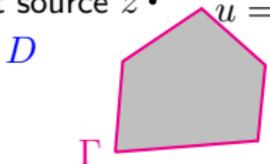
$$u(x) = u^i(x) - \int_{\Gamma} \Phi(x, y) \partial_n u(y) \, ds(y), \quad x \in D.$$

Taking a linear combination of Dirichlet ( $\gamma$ ) and Neumann ( $\partial_n$ ) traces, we obtain the standard 2nd kind integral equation

$$A \partial_n u = g := \partial_n u^i - ik \gamma u^i, \quad \text{where} \quad A := \frac{1}{2} I + K' - ikS.$$

## Coercive formulations: bounded obstacle case

Point source  $z \bullet$



The diagram shows a gray pentagonal domain  $D$  with a pink boundary  $\Gamma$ . A point source  $z$  is marked with a black dot inside the domain. The boundary is labeled  $u = 0$  at the top vertex.

$$\Delta u + k^2 u = \delta_z, \quad k > 0$$

$$u(x) = u^i(x) - \int_{\Gamma} \Phi(x, y) \partial_n u(y) \, ds(y), \quad x \in D.$$

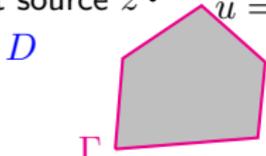
Taking a linear combination of Dirichlet ( $\gamma$ ) and Neumann ( $\partial_n$ ) and surface gradient ( $\nabla_{\Gamma}$ ) traces, we obtain the 2nd kind integral equation

$$A_Z \partial_n u = g_Z := Z \cdot \gamma \nabla u^i - ik \gamma u^i, \quad \text{where} \quad A_Z := Z \cdot n \left( \frac{1}{2} I + K' \right) + Z \cdot \nabla_{\Gamma} S - i \alpha S,$$

and  $Z : \Gamma \rightarrow \mathbb{R}^d$  is in  $L^{\infty}(\Gamma)$ .

## Coercive formulations: bounded obstacle case

Point source  $z \bullet$



The diagram shows a gray shaded pentagonal domain  $D$  with a pink boundary  $\Gamma$ . A point source  $z$  is marked with a black dot inside the domain. The boundary  $\Gamma$  is labeled with  $u = 0$  at the top vertex.

$$\Delta u + k^2 u = \delta_z, \quad k > 0$$

$$u(x) = u^i(x) - \int_{\Gamma} \Phi(x, y) \partial_n u(y) \, ds(y), \quad x \in D.$$

Taking a linear combination of Dirichlet ( $\gamma$ ) and Neumann ( $\partial_n$ ) and surface gradient ( $\nabla_{\Gamma}$ ) traces, we obtain the 2nd kind integral equation

$$A_Z \partial_n u = g_Z := Z \cdot \gamma \nabla u^i - ik \gamma u^i, \quad \text{where} \quad A_Z := Z \cdot n \left( \frac{1}{2} I + K' \right) + Z \cdot \nabla_{\Gamma} S - i \alpha S,$$

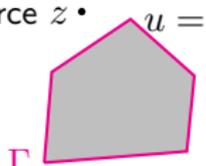
and  $Z : \Gamma \rightarrow \mathbb{R}^d$  is in  $L^{\infty}(\Gamma)$ . If  $Z = n$  and  $\alpha = k$ , then  $A_Z = A$  and  $g_Z = g$ .

## Coercive formulations: bounded obstacle case

Point source  $z \bullet$

$D$

$u = 0$



$\Delta u + k^2 u = \delta_z, \quad k > 0$

$$u(x) = u^i(x) - \int_{\Gamma} \Phi(x, y) \partial_n u(y) \, ds(y), \quad x \in D.$$

Taking a linear combination of Dirichlet ( $\gamma$ ) and Neumann ( $\partial_n$ ) and surface gradient ( $\nabla_{\Gamma}$ ) traces, we obtain the 2nd kind integral equation

$$A_Z \partial_n u = g_Z := Z \cdot \gamma \nabla u^i - ik \gamma u^i, \quad \text{where} \quad A_Z := Z \cdot n \left( \frac{1}{2} I + K' \right) + Z \cdot \nabla_{\Gamma} S - i \alpha S,$$

and  $Z : \Gamma \rightarrow \mathbb{R}^d$  is in  $L^{\infty}(\Gamma)$ . If  $Z = n$  and  $\alpha = k$ , then  $A_Z = A$  and  $g_Z = g$ .

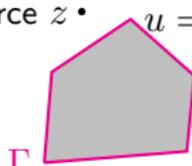
**Theorem** (C-W & Spence 2022b). If  $Z$  is continuous and  $Z \cdot n \geq c > 0$  on  $\Gamma$ , then  $A_Z = A_0 + K$  where  $A_0$  is coercive and  $K$  is compact, so that all Galerkin methods for  $A_Z \partial_n u = g_Z$  are convergent, provided  $A_Z$  is injective.

## Coercive formulations: bounded obstacle case

Point source  $z \bullet$

$D$

$u = 0$



$\Delta u + k^2 u = \delta_z, \quad k > 0$

$$u(x) = u^i(x) - \int_{\Gamma} \Phi(x, y) \partial_n u(y) ds(y), \quad x \in D.$$

Taking a linear combination of Dirichlet ( $\gamma$ ) and Neumann ( $\partial_n$ ) and surface gradient ( $\nabla_{\Gamma}$ ) traces, we obtain the 2nd kind integral equation

$$A_Z \partial_n u = g_Z := Z \cdot \gamma \nabla u^i - ik \gamma u^i, \quad \text{where} \quad A_Z := Z \cdot n \left( \frac{1}{2} I + K' \right) + Z \cdot \nabla_{\Gamma} S - i \alpha S,$$

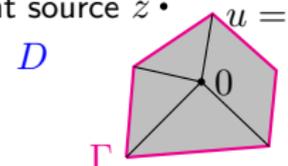
and  $Z : \Gamma \rightarrow \mathbb{R}^d$  is in  $L^{\infty}(\Gamma)$ . If  $Z = n$  and  $\alpha = k$ , then  $A_Z = A$  and  $g_Z = g$ .

**Theorem** (C-W & Spence 2022b). If  $Z$  is continuous and  $Z \cdot n \geq c > 0$  on  $\Gamma$ , then  $A_Z = A_0 + K$  where  $A_0$  is coercive and  $K$  is compact, so that all Galerkin methods for  $A_Z \partial_n u = g_Z$  are convergent, provided  $A_Z$  is injective.

Sadly injectivity of  $A_Z$  not yet clear in general (**open problem**).

# Coercive formulations: bounded obstacle case

Point source  $z \bullet$



$D$

$\Gamma$

$u = 0$

$0$

$\Delta u + k^2 u = \delta_z, \quad k > 0$

$$u(x) = u^i(x) - \int_{\Gamma} \Phi(x, y) \partial_n u(y) \, ds(y), \quad x \in D.$$

Taking a linear combination of Dirichlet ( $\gamma$ ) and Neumann ( $\partial_n$ ) and surface gradient ( $\nabla_{\Gamma}$ ) traces, we obtain the 2nd kind integral equation

$$A_Z \partial_n u = g_Z := Z \cdot \gamma \nabla u^i - ik \gamma u^i, \quad \text{where} \quad A_Z := Z \cdot n \left( \frac{1}{2} I + K' \right) + Z \cdot \nabla_{\Gamma} S - i \alpha S,$$

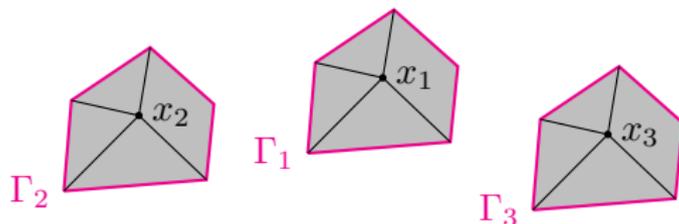
and  $Z : \Gamma \rightarrow \mathbb{R}^d$  is in  $L^{\infty}(\Gamma)$ . If  $Z = n$  and  $\alpha = k$ , then  $A_Z = A$  and  $g_Z = g$ .

**Theorem** (Spence, C-W, Graham, Smyshlyaev 2011). If  $\Gamma$  is **star-shaped** with respect to 0,

$$Z(x) := x, \quad \alpha(x) := k|x| + i(d-1)/2, \quad x \cdot n \geq c > 0,$$

on  $\Gamma$ , then  $A_Z$  is **uniformly-in- $k$  coercive** with coercivity constant  $\gamma = c/2$ , so that all Galerkin methods for  $A_Z \partial_n u = g_Z$  are convergent.

# Multiple scattering formulation



$$u(x) = u^i(x) - \int_{\Gamma} \Phi(x, y) \partial_n u(y) \, ds(y), \quad x \in D.$$

$$A_Z \partial_n u = g_Z := Z \cdot \gamma \nabla u^i - ik \gamma u^i, \quad \text{where} \quad A_Z := Z \cdot n \left( \frac{1}{2} I + K' \right) + Z \cdot \nabla_{\Gamma} S - i \alpha S.$$

**Corollary** (Gibbs, C-W, Langdon, Moiola 2021). If each component  $\Gamma_j$  of  $\Gamma$  is star-shaped, and, on  $\Gamma_j$ ,

$$Z(x) := x - x_j, \quad \alpha(x) := k|x - x_j| + i(d-1)/2, \quad (x - x_j) \cdot n \geq c > 0,$$

then  $A_Z = A_0 + K$  with  $A_0$  coercive and  $K$  compact, **and**  $A_Z$  is injective, so that all Galerkin methods for  $A_Z \partial_n u = g_Z$  are convergent.

## Our typical RSS problem

Suppose  $d = 2$  or  $3$ ,  $f : \mathbb{R}^{d-1} \rightarrow \mathbb{R}$  is bounded and Lipschitz continuous, precisely

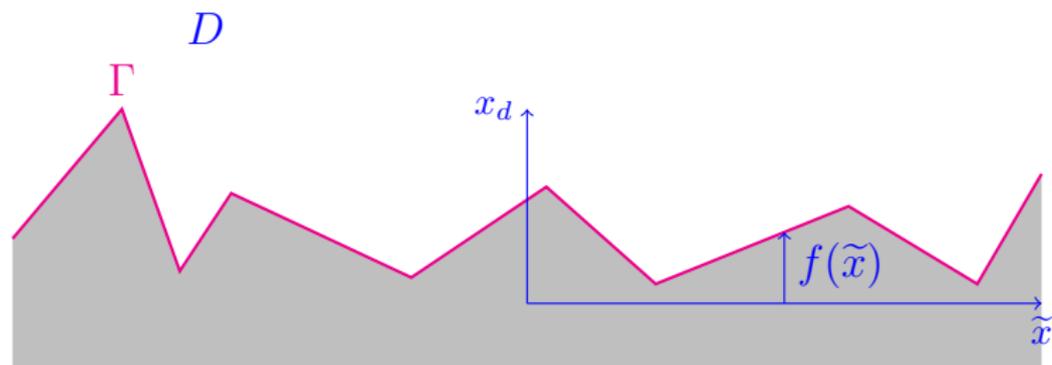
$$0 < f_- \leq f(\tilde{x}) \leq f_+ \quad \text{and} \quad |f(\tilde{x}) - f(\tilde{y})| \leq L|\tilde{x} - \tilde{y}|, \quad \tilde{x}, \tilde{y} \in \mathbb{R}^{d-1}.$$

## Our typical RSS problem; the rough surface is $\Gamma$

Suppose  $d = 2$  or  $3$ ,  $f : \mathbb{R}^{d-1} \rightarrow \mathbb{R}$  is bounded and Lipschitz continuous, precisely

$$0 < f_- \leq f(\tilde{x}) \leq f_+ \quad \text{and} \quad |f(\tilde{x}) - f(\tilde{y})| \leq L|\tilde{x} - \tilde{y}|, \quad \tilde{x}, \tilde{y} \in \mathbb{R}^{d-1}.$$

Let



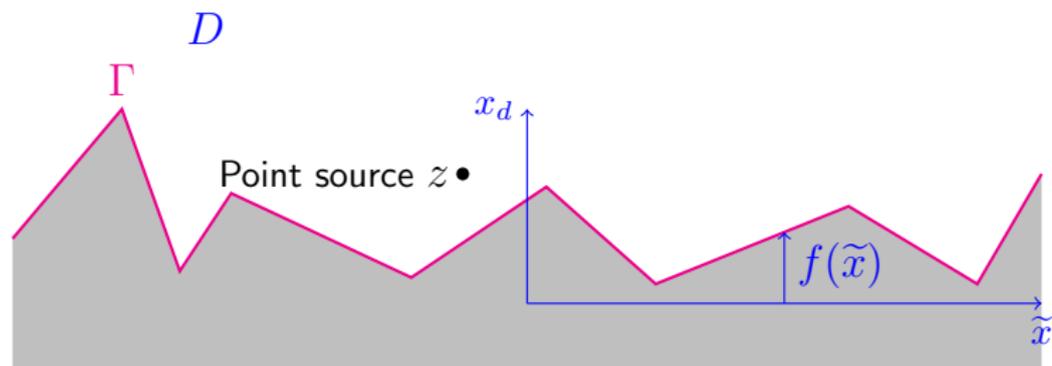
$$D := \{(\tilde{x}, x_d) : x_d > f(\tilde{x}), \tilde{x} \in \mathbb{R}^{d-1}\} \subset \mathbb{R}^d, \quad \Gamma := \partial D = \{(\tilde{x}, f(\tilde{x})) : \tilde{x} \in \mathbb{R}^{d-1}\}.$$

## Our typical RSS problem; the rough surface is $\Gamma$

Suppose  $d = 2$  or  $3$ ,  $f : \mathbb{R}^{d-1} \rightarrow \mathbb{R}$  is bounded and Lipschitz continuous, precisely

$$0 < f_- \leq f(\tilde{x}) \leq f_+ \quad \text{and} \quad |f(\tilde{x}) - f(\tilde{y})| \leq L|\tilde{x} - \tilde{y}|, \quad \tilde{x}, \tilde{y} \in \mathbb{R}^{d-1}.$$

Let



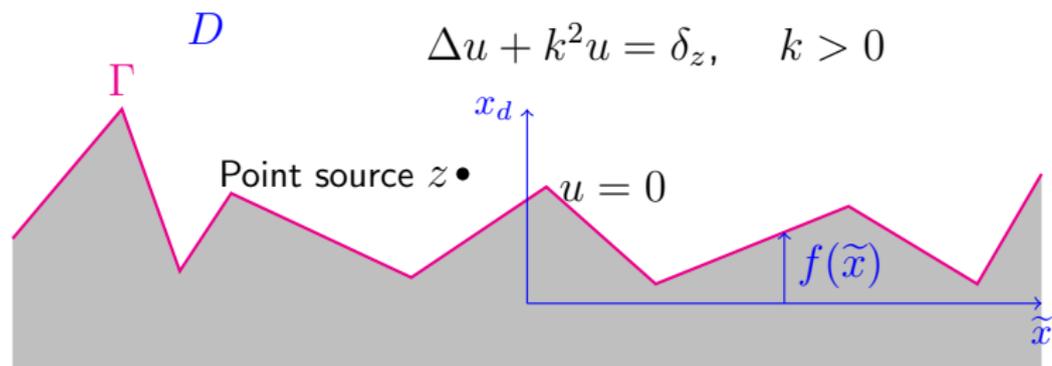
$$D := \{(\tilde{x}, x_d) : x_d > f(\tilde{x}), \tilde{x} \in \mathbb{R}^{d-1}\} \subset \mathbb{R}^d, \quad \Gamma := \partial D = \{(\tilde{x}, f(\tilde{x})) : \tilde{x} \in \mathbb{R}^{d-1}\}.$$

## Our typical RSS problem; the rough surface is $\Gamma$

Suppose  $d = 2$  or  $3$ ,  $f : \mathbb{R}^{d-1} \rightarrow \mathbb{R}$  is bounded and Lipschitz continuous, precisely

$$0 < f_- \leq f(\tilde{x}) \leq f_+ \quad \text{and} \quad |f(\tilde{x}) - f(\tilde{y})| \leq L|\tilde{x} - \tilde{y}|, \quad \tilde{x}, \tilde{y} \in \mathbb{R}^{d-1}.$$

Let



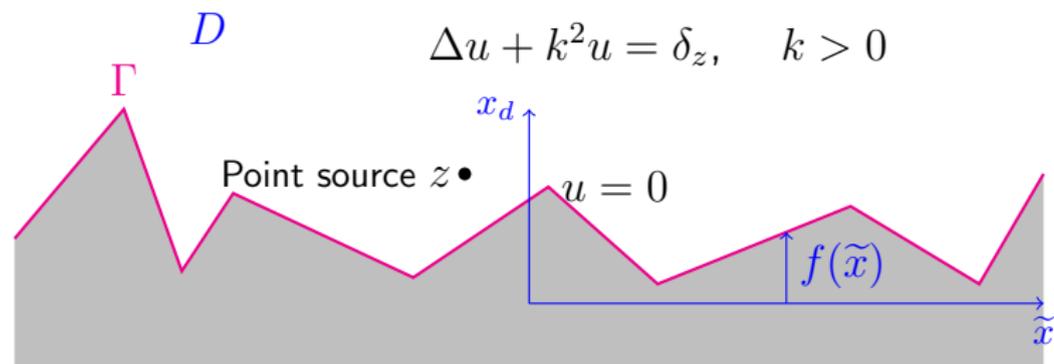
$$D := \{(\tilde{x}, x_d) : x_d > f(\tilde{x}), \tilde{x} \in \mathbb{R}^{d-1}\} \subset \mathbb{R}^d, \quad \Gamma := \partial D = \{(\tilde{x}, f(\tilde{x})) : \tilde{x} \in \mathbb{R}^{d-1}\}.$$

# Our typical RSS problem; the rough surface is $\Gamma$

Suppose  $d = 2$  or  $3$ ,  $f : \mathbb{R}^{d-1} \rightarrow \mathbb{R}$  is bounded and Lipschitz continuous, precisely

$$0 < f_- \leq f(\tilde{x}) \leq f_+ \quad \text{and} \quad |f(\tilde{x}) - f(\tilde{y})| \leq L|\tilde{x} - \tilde{y}|, \quad \tilde{x}, \tilde{y} \in \mathbb{R}^{d-1}.$$

Let



$$D := \{(\tilde{x}, x_d) : x_d > f(\tilde{x}), \tilde{x} \in \mathbb{R}^{d-1}\} \subset \mathbb{R}^d, \quad \Gamma := \partial D = \{(\tilde{x}, f(\tilde{x})) : \tilde{x} \in \mathbb{R}^{d-1}\}.$$

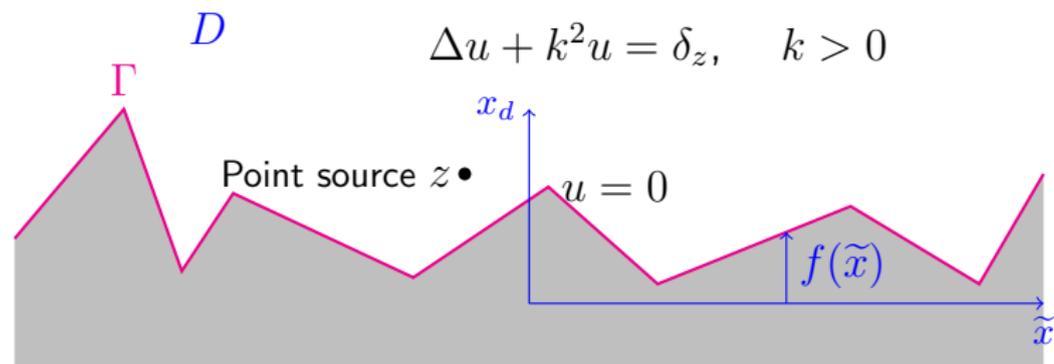
**Key feature:**  $\Gamma$  unbounded (in the horizontal directions).

# Our typical RSS problem; the rough surface is $\Gamma$

Suppose  $d = 2$  or  $3$ ,  $f : \mathbb{R}^{d-1} \rightarrow \mathbb{R}$  is bounded and Lipschitz continuous, precisely

$$0 < f_- \leq f(\tilde{x}) \leq f_+ \quad \text{and} \quad |f(\tilde{x}) - f(\tilde{y})| \leq L|\tilde{x} - \tilde{y}|, \quad \tilde{x}, \tilde{y} \in \mathbb{R}^{d-1}.$$

Let



$$D := \{(\tilde{x}, x_d) : x_d > f(\tilde{x}), \tilde{x} \in \mathbb{R}^{d-1}\} \subset \mathbb{R}^d, \quad \Gamma := \partial D = \{(\tilde{x}, f(\tilde{x})) : \tilde{x} \in \mathbb{R}^{d-1}\}.$$

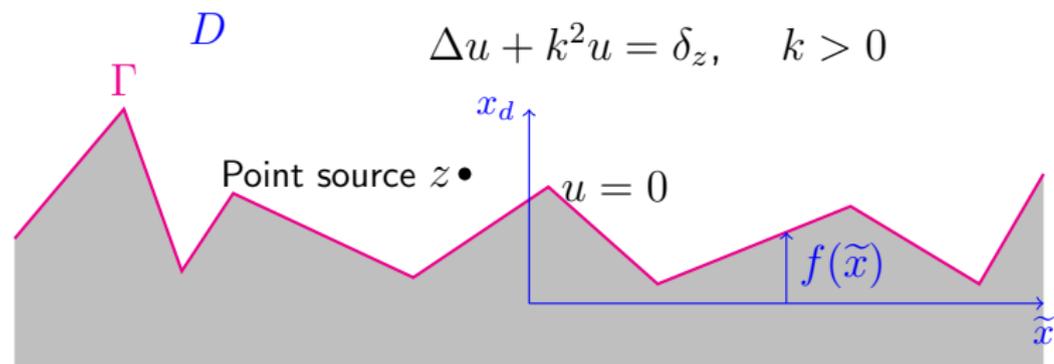
**Key feature:**  $\Gamma$  **unbounded** (in the horizontal directions). The **dimensionless surface elevation**,  $k(f_+ - f_-)$ , need not be large.

# Our typical RSS problem; the rough surface is $\Gamma$

Suppose  $d = 2$  or  $3$ ,  $f : \mathbb{R}^{d-1} \rightarrow \mathbb{R}$  is bounded and Lipschitz continuous, precisely

$$0 < f_- \leq f(\tilde{x}) \leq f_+ \quad \text{and} \quad |f(\tilde{x}) - f(\tilde{y})| \leq L|\tilde{x} - \tilde{y}|, \quad \tilde{x}, \tilde{y} \in \mathbb{R}^{d-1}.$$

Let

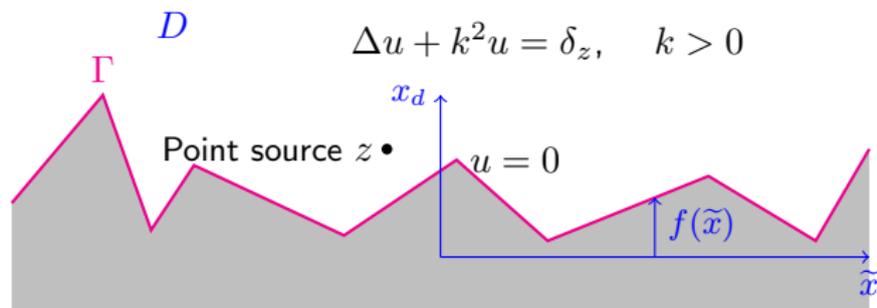


$$D := \{(\tilde{x}, x_d) : x_d > f(\tilde{x}), \tilde{x} \in \mathbb{R}^{d-1}\} \subset \mathbb{R}^d, \quad \Gamma := \partial D = \{(\tilde{x}, f(\tilde{x})) : \tilde{x} \in \mathbb{R}^{d-1}\}.$$

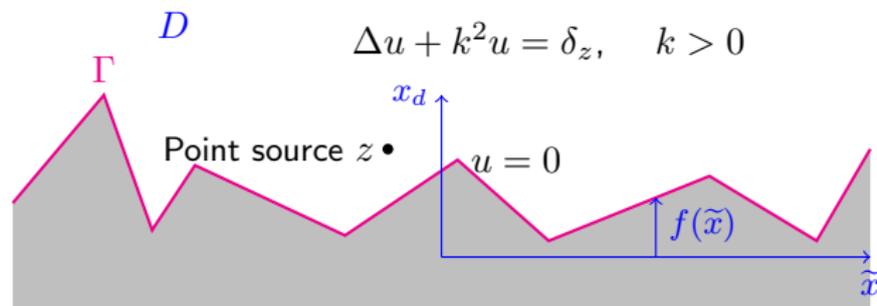
**Key feature:**  $\Gamma$  **unbounded** (in the horizontal directions). The **dimensionless surface elevation**,  $k(f_+ - f_-)$ , need not be large.

Applications in outdoor noise or radar propagation over ground and sea surfaces, and in optics: all nominally flat surfaces are rough at some scale!

# Integral equation methods: rough surface scattering



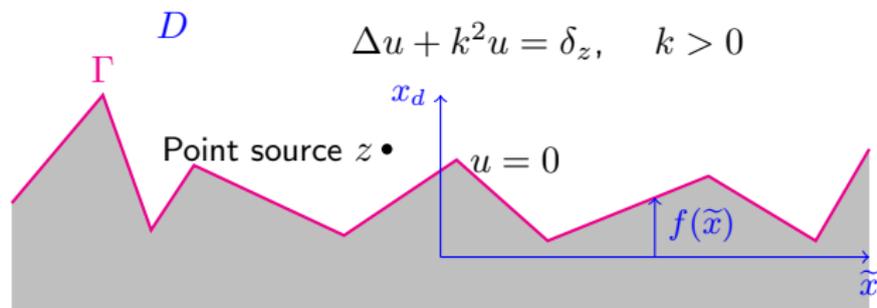
# Integral equation methods: rough surface scattering



**First idea:** just use the **bounded obstacle formulation**, i.e.

$$u(x) = u^i(x) - \int_{\Gamma} \Phi(x, y) \partial_n u(y) \, ds(y), \quad x \in D.$$

# Integral equation methods: rough surface scattering



**First idea:** just use the **bounded obstacle formulation**, i.e.

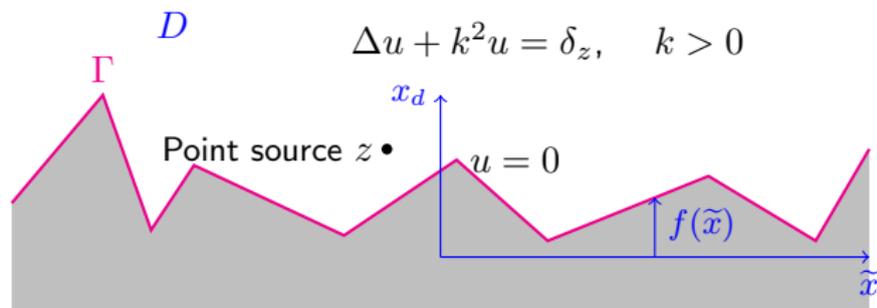
$$u(x) = u^i(x) - \int_{\Gamma} \Phi(x, y) \partial_n u(y) ds(y), \quad x \in D.$$

where  $\partial_n u$  satisfies

$$A \partial_n u = g := \partial_n u - ik\gamma u, \quad \text{and} \quad A := \frac{1}{2}I + K' - ikS,$$

$$K' \varphi(x) := \int_{\Gamma} \partial_{n(x)} \Phi(x, y) \varphi(y) ds(y), \quad S \varphi(x) := \int_{\Gamma} \Phi(x, y) \varphi(y) ds(y), \quad x \in \Gamma.$$

# Integral equation methods: rough surface scattering



**First idea:** just use the **bounded obstacle formulation**, i.e.

$$u(x) = u^i(x) - \int_{\Gamma} \Phi(x, y) \partial_n u(y) ds(y), \quad x \in D.$$

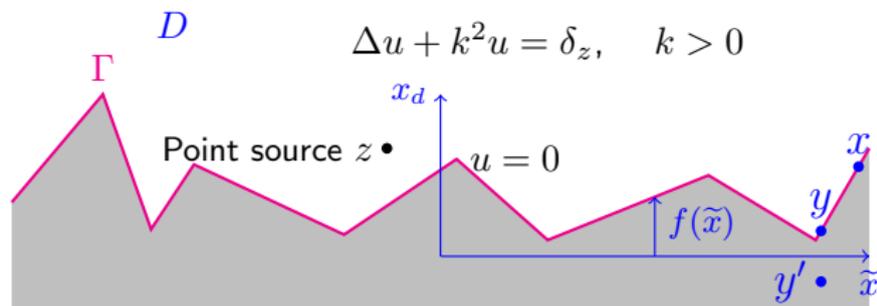
where  $\partial_n u$  satisfies

$$A \partial_n u = g := \partial_n u - ik\gamma u, \quad \text{and} \quad A := \frac{1}{2}I + K' - ikS,$$

$$K' \varphi(x) := \int_{\Gamma} \partial_{n(x)} \Phi(x, y) \varphi(y) ds(y), \quad S \varphi(x) := \int_{\Gamma} \Phi(x, y) \varphi(y) ds(y), \quad x \in \Gamma.$$

**Issue:**  $\Phi(x, y)$  decays too slowly for  $A$  to be a bounded operator.

# Integral equation methods: rough surface scattering



**First idea:** just use the **bounded obstacle formulation**, i.e.

$$u(x) = u^i(x) - \int_{\Gamma} G(x, y) \partial_n u(y) ds(y), \quad x \in D.$$

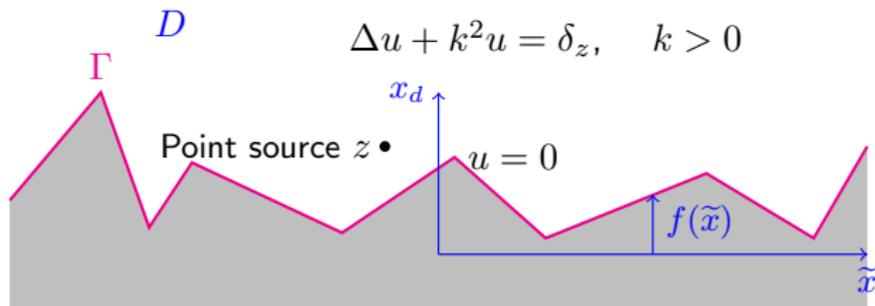
where  $\partial_n u$  satisfies

$$A \partial_n u = g := \partial_n u - ik\gamma u, \quad \text{and} \quad A := \frac{1}{2}I + K' - ikS,$$

$$K' \varphi(x) := \int_{\Gamma} \partial_{n(x)} G(x, y) \varphi(y) ds(y), \quad S \varphi(x) := \int_{\Gamma} G(x, y) \varphi(y) ds(y), \quad x \in \Gamma.$$

**Issue:**  $\Phi(x, y)$  decays too slowly for  $A$  to be a bounded operator.

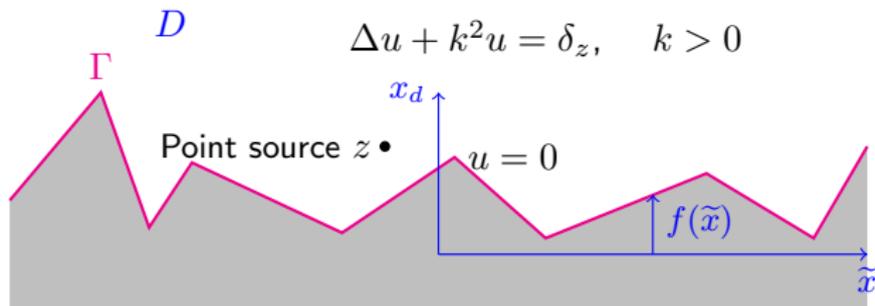
**Solution:** (Zhang & C-W 2003, C-W, Heinemeyer, Potthast 2006a,b) Replace  $\Phi(x, y)$  with Dirichlet half-space Green's function,  $G(x, y) := \Phi(x, y) - \Phi(x, y')$ .



$$A\partial_n u = g := \partial_n u - ik\gamma u, \quad \text{and} \quad A := \frac{1}{2}I + K' - ikS,$$

$$K'\varphi(x) := \int_{\Gamma} \partial_{n(x)} G(x, y) \varphi(y) ds(y), \quad S\varphi(x) := \int_{\Gamma} G(x, y) \varphi(y) ds(y), \quad x \in \Gamma.$$

**Solution:** (Zhang & C-W 2003, C-W, Heinemeyer, Potthast 2006a,b) Replace  $\Phi(x, y)$  with Dirichlet half-space Green's function,  $G(x, y) := \Phi(x, y) - \Phi(x, y')$ .



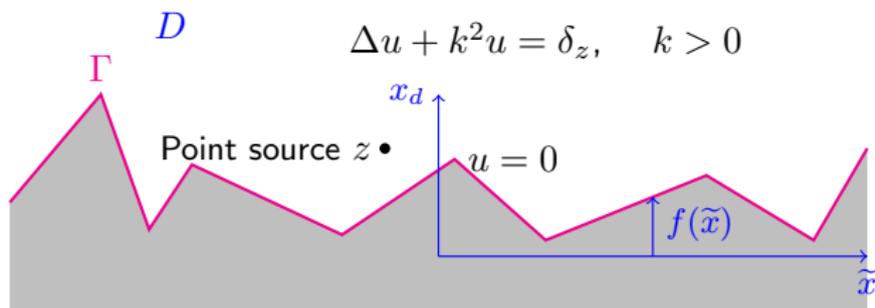
$$A\partial_n u = g := \partial_n u - ik\gamma u, \quad \text{and} \quad A := \frac{1}{2}I + K' - ikS,$$

$$K'\varphi(x) := \int_{\Gamma} \partial_{n(x)} G(x, y) \varphi(y) ds(y), \quad S\varphi(x) := \int_{\Gamma} G(x, y) \varphi(y) ds(y), \quad x \in \Gamma.$$

**Solution:** (Zhang & C-W 2003, C-W, Heinemeyer, Potthast 2006a,b) Replace  $\Phi(x, y)$  with Dirichlet half-space Green's function,  $G(x, y) := \Phi(x, y) - \Phi(x, y')$ .

**Theorem** (C-W, Heinemeyer, Potthast 2006a,b).  $A$  is bounded and invertible on  $L^2(\Gamma)$ , indeed, where  $L$  is the Lipschitz constant of  $f$ ,

$$\|A^{-1}\| \leq 12(1 + L)^2.$$



$$A \partial_n u = g := \partial_n u - ik\gamma u, \quad \text{and} \quad A := \frac{1}{2}I + K' - ikS,$$

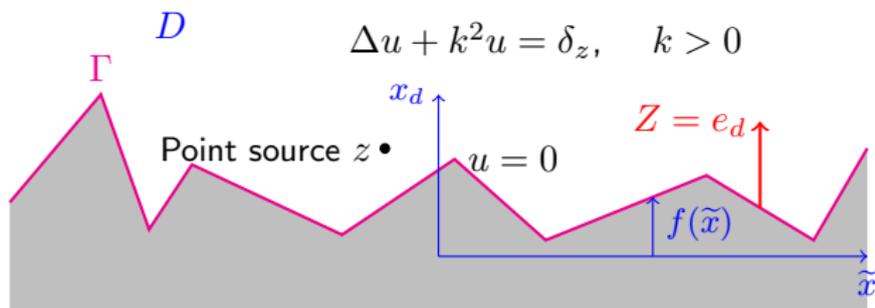
$$K' \varphi(x) := \int_{\Gamma} \partial_{n(x)} G(x, y) \varphi(y) ds(y), \quad S \varphi(x) := \int_{\Gamma} G(x, y) \varphi(y) ds(y), \quad x \in \Gamma.$$

**Solution:** (Zhang & C-W 2003, C-W, Heinemeyer, Potthast 2006a,b) Replace  $\Phi(x, y)$  with Dirichlet half-space Green's function,  $G(x, y) := \Phi(x, y) - \Phi(x, y')$ .

**Theorem** (C-W, Heinemeyer, Potthast 2006a,b).  $A$  is bounded and invertible on  $L^2(\Gamma)$ , indeed, where  $L$  is the Lipschitz constant of  $f$ ,

$$\|A^{-1}\| \leq 12(1 + L)^2.$$

**Issue:** but how do we prove convergence of boundary truncation, BEM, GMRES?



$$A_Z \partial_n u = g_Z := Z \cdot \gamma \nabla u^i - ik \gamma u^i, \quad \text{where} \quad A_Z := Z \cdot n \left( \frac{1}{2} I + K' \right) + Z \cdot \nabla_{\Gamma} S - ik S.$$

$$K' \varphi(x) := \int_{\Gamma} \partial_{n(x)} G(x, y) \varphi(y) ds(y), \quad S \varphi(x) := \int_{\Gamma} G(x, y) \varphi(y) ds(y), \quad x \in \Gamma.$$

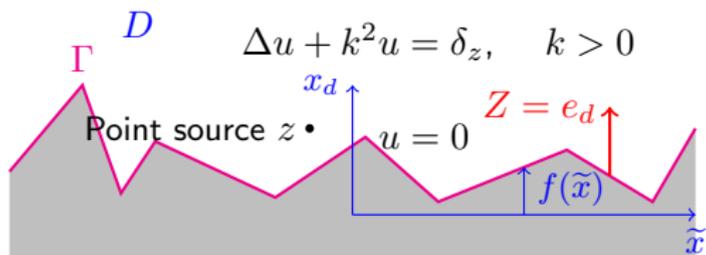
**Solution:** (Zhang & C-W 2003, C-W, Heinemeyer, Potthast 2006a,b) Replace  $\Phi(x, y)$  with Dirichlet half-space Green's function,  $G(x, y) := \Phi(x, y) - \Phi(x, y')$ .

**Theorem** (C-W, Heinemeyer, Potthast 2006a,b).  $A$  is bounded and invertible on  $L^2(\Gamma)$ , indeed, where  $L$  is the Lipschitz constant of  $f$ ,

$$\|A^{-1}\| \leq 12(1 + L)^2.$$

**Issue:** but how do we prove convergence of boundary truncation, BEM, GMRES?

**Solution:** replace  $A$  with  $A_Z$  with  $Z = e_d$ , so that  $Z \cdot n \geq (1 + L^2)^{-1/2}$ .



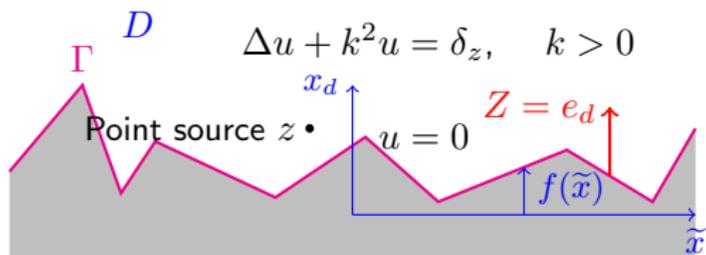
$$A_Z \partial_n u = g_Z := Z \cdot \gamma \nabla u^i - ik \gamma u^i, \quad \text{where} \quad A_Z := Z \cdot n \left( \frac{1}{2} I + K' \right) + Z \cdot \nabla_{\Gamma} S - ik S.$$

**Theorem** (C-W, Heinemeyer, Potthast 2006a,b).  $A$  is bounded and invertible on  $L^2(\Gamma)$ , indeed, where  $L$  is the Lipschitz constant of  $f$ ,

$$\|A^{-1}\| \leq 12(1 + L)^2.$$

**Theorem.** With  $Z = e_d$ ,  $A_Z$  is bounded and **uniformly-in- $k$  coercive** on  $L^2(\Gamma)$ , with coercivity constant

$$\gamma := \frac{1}{2(1 + L^2)^{1/2}} \quad \text{so that} \quad \|A_Z^{-1}\| \leq \gamma^{-1} = 2(1 + L^2)^{1/2}.$$



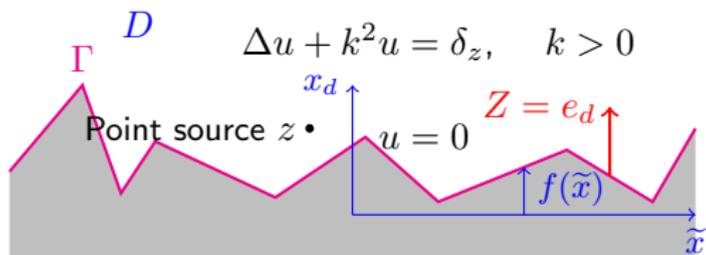
$$A_Z \partial_n u = g_Z := Z \cdot \gamma \nabla u^i - ik \gamma u^i, \quad \text{where} \quad A_Z := Z \cdot n \left( \frac{1}{2} I + K' \right) + Z \cdot \nabla_{\Gamma} S - ik S.$$

**Theorem.** With  $Z = e_d$ ,  $A_Z$  is bounded and **uniformly-in- $k$  coercive** on  $L^2(\Gamma)$ , with coercivity constant

$$\gamma := \frac{1}{2(1 + L^2)^{1/2}} \quad \text{so that} \quad \|A_Z^{-1}\| \leq \gamma^{-1} = 2(1 + L^2)^{1/2}.$$

Thus, if  $H_N \subset L^2(\Gamma)$  is any BEM subspace supported on a finite part of  $\Gamma$  of diameter  $2a$ , then the Galerkin approximation  $\varphi_N \in H_N$  to  $\partial_n u$  is well-defined and

$$\|\partial_n u - \varphi_N\| \leq \frac{\|A_Z\|}{\gamma} \inf_{\psi_N \in H_N} \|\partial_n u - \psi_N\|.$$



$$A_Z \partial_n u = g_Z := Z \cdot \gamma \nabla u^i - ik \gamma u^i, \quad \text{where} \quad A_Z := Z \cdot n \left( \frac{1}{2} I + K' \right) + Z \cdot \nabla_{\Gamma} S - ik S.$$

**Theorem.** With  $Z = e_d$ ,  $A_Z$  is bounded and **uniformly-in- $k$  coercive** on  $L^2(\Gamma)$ , with coercivity constant

$$\gamma := \frac{1}{2(1 + L^2)^{1/2}} \quad \text{so that} \quad \|A_Z^{-1}\| \leq \gamma^{-1} = 2(1 + L^2)^{1/2}.$$

Thus, if  $H_N \subset L^2(\Gamma)$  is any BEM subspace supported on a finite part of  $\Gamma$  of diameter  $2a$ , then the Galerkin approximation  $\varphi_N \in H_N$  to  $\partial_n u$  is well-defined and

$$\|\partial_n u - \varphi_N\| \leq \frac{\|A_Z\|}{\gamma} \inf_{\psi_N \in H_N} \|\partial_n u - \psi_N\|.$$

Moreover, if  $M$  is the mass matrix of the chosen basis for  $H_N$  and  $r_m$  is the residual after  $m$  steps of GMRES,

$$\frac{\|r_m\|_2}{\|r_0\|_2} \leq \varepsilon \quad \text{provided} \quad m \geq \frac{3\sqrt{3}}{4} \frac{\|A_Z\|}{\gamma} \text{cond}(M) \log \left( \frac{8}{\varepsilon} \right).$$

**Theorem.** With  $Z = e_d$ ,  $A_Z$  is bounded and **uniformly-in- $k$  coercive** on  $L^2(\Gamma)$ , with coercivity constant

$$\gamma := \frac{1}{2(1+L^2)^{1/2}} \quad \text{so that} \quad \|A_Z^{-1}\| \leq \gamma^{-1} = 2(1+L^2)^{1/2}.$$

Thus, if  $H_N \subset L^2(\Gamma)$  is any BEM subspace supported on a finite part of  $\Gamma$  of diameter  $2a$ , then the Galerkin approximation  $\varphi_N \in H_N$  to  $\partial_n u$  is well-defined and

$$\|\partial_n u - \varphi_N\| \leq \frac{\|A_Z\|}{\gamma} \inf_{\psi_N \in H_N} \|\partial_n u - \psi_N\|.$$

Moreover, if  $M$  is the mass matrix of the chosen basis for  $H_N$  and  $r_m$  is the residual after  $m$  steps of GMRES,

$$\frac{\|r_m\|_2}{\|r_0\|_2} \leq \varepsilon \quad \text{provided} \quad m \geq \frac{3\sqrt{3}}{4} \frac{\|A_Z\|}{\gamma} \text{cond}(M) \log\left(\frac{8}{\varepsilon}\right).$$

*Idea of proof.* The proof combines:

- harmonic analysis techniques for 2nd kind integral equations on Lipschitz domains

**Theorem.** With  $Z = e_d$ ,  $A_Z$  is bounded and **uniformly-in- $k$  coercive** on  $L^2(\Gamma)$ , with coercivity constant

$$\gamma := \frac{1}{2(1+L^2)^{1/2}} \quad \text{so that} \quad \|A_Z^{-1}\| \leq \gamma^{-1} = 2(1+L^2)^{1/2}.$$

Thus, if  $H_N \subset L^2(\Gamma)$  is any BEM subspace supported on a finite part of  $\Gamma$  of diameter  $2a$ , then the Galerkin approximation  $\varphi_N \in H_N$  to  $\partial_n u$  is well-defined and

$$\|\partial_n u - \varphi_N\| \leq \frac{\|A_Z\|}{\gamma} \inf_{\psi_N \in H_N} \|\partial_n u - \psi_N\|.$$

Moreover, if  $M$  is the mass matrix of the chosen basis for  $H_N$  and  $r_m$  is the residual after  $m$  steps of GMRES,

$$\frac{\|r_m\|_2}{\|r_0\|_2} \leq \varepsilon \quad \text{provided} \quad m \geq \frac{3\sqrt{3}}{4} \frac{\|A_Z\|}{\gamma} \text{cond}(M) \log\left(\frac{8}{\varepsilon}\right).$$

*Idea of proof.* The proof combines:

- harmonic analysis techniques for 2nd kind integral equations on Lipschitz domains
- methods for proving invertibility/coercivity through Rellich-type identities, combining ideas of Verchota (1984), C-W and Monk (2005), C-W, Heinemeyer, Potthast (2006b), Spence, C-W, Graham, Smyshlyaev (2011).

**Theorem.** With  $Z = e_d$ ,  $A_Z$  is bounded and **uniformly-in- $k$  coercive** on  $L^2(\Gamma)$ , with coercivity constant

$$\gamma := \frac{1}{2(1+L^2)^{1/2}} \quad \text{so that} \quad \|A_Z^{-1}\| \leq \gamma^{-1} = 2(1+L^2)^{1/2}.$$

Thus, if  $H_N \subset L^2(\Gamma)$  is any BEM subspace supported on a finite part of  $\Gamma$  of diameter  $2a$ , then the Galerkin approximation  $\varphi_N \in H_N$  to  $\partial_n u$  is well-defined and

$$\|\partial_n u - \varphi_N\| \leq \frac{\|A_Z\|}{\gamma} \inf_{\psi_N \in H_N} \|\partial_n u - \psi_N\|.$$

Moreover, if  $M$  is the mass matrix of the chosen basis for  $H_N$  and  $r_m$  is the residual after  $m$  steps of GMRES,

$$\frac{\|r_m\|_2}{\|r_0\|_2} \leq \varepsilon \quad \text{provided} \quad m \geq \frac{3\sqrt{3}}{4} \frac{\|A_Z\|}{\gamma} \text{cond}(M) \log\left(\frac{8}{\varepsilon}\right).$$

*Idea of proof.* The proof combines:

- harmonic analysis techniques for 2nd kind integral equations on Lipschitz domains
- methods for proving invertibility/coercivity through Rellich-type identities, combining ideas of Verchota (1984), C-W and Monk (2005), C-W, Heinemeyer, Potthast (2006b), Spence, C-W, Graham, Smyshlyaev (2011). The Rellich identity we need follows from writing in divergence form integrals of the form

$$\int (\Delta u + k^2 u) \frac{\partial \bar{u}}{\partial x_d} dx.$$

**Theorem.** With  $Z = e_d$ ,  $A_Z$  is bounded and **uniformly-in- $k$  coercive** on  $L^2(\Gamma)$ , with coercivity constant

$$\gamma := \frac{1}{2(1+L^2)^{1/2}} \quad \text{so that} \quad \|A_Z^{-1}\| \leq \gamma^{-1} = 2(1+L^2)^{1/2}.$$

Thus, if  $H_N \subset L^2(\Gamma)$  is any BEM subspace supported on a finite part of  $\Gamma$  of diameter  $2a$ , then the Galerkin approximation  $\varphi_N \in H_N$  to  $\partial_n u$  is well-defined and

$$\|\partial_n u - \varphi_N\| \leq \frac{\|A_Z\|}{\gamma} \inf_{\psi_N \in H_N} \|\partial_n u - \psi_N\|.$$

Moreover, if  $M$  is the mass matrix of the chosen basis for  $H_N$  and  $r_m$  is the residual after  $m$  steps of GMRES,

$$\frac{\|r_m\|_2}{\|r_0\|_2} \leq \varepsilon \quad \text{provided} \quad m \geq \frac{3\sqrt{3}}{4} \frac{\|A_Z\|}{\gamma} \text{cond}(M) \log\left(\frac{8}{\varepsilon}\right).$$

*Idea of proof.* The proof combines:

- harmonic analysis techniques for 2nd kind integral equations on Lipschitz domains
- methods for proving invertibility/coercivity through Rellich-type identities, combining ideas of Verchota (1984), C-W and Monk (2005), C-W, Heinemeyer, Potthast (2006b), Spence, C-W, Graham, Smyshlyaev (2011). The Rellich identity we need follows from writing in divergence form integrals of the form

$$\int (\Delta u + k^2 u) \frac{\partial \bar{u}}{\partial x_d} dx.$$

- The convergence theory for Galerkin BEM and GMRES recalled earlier

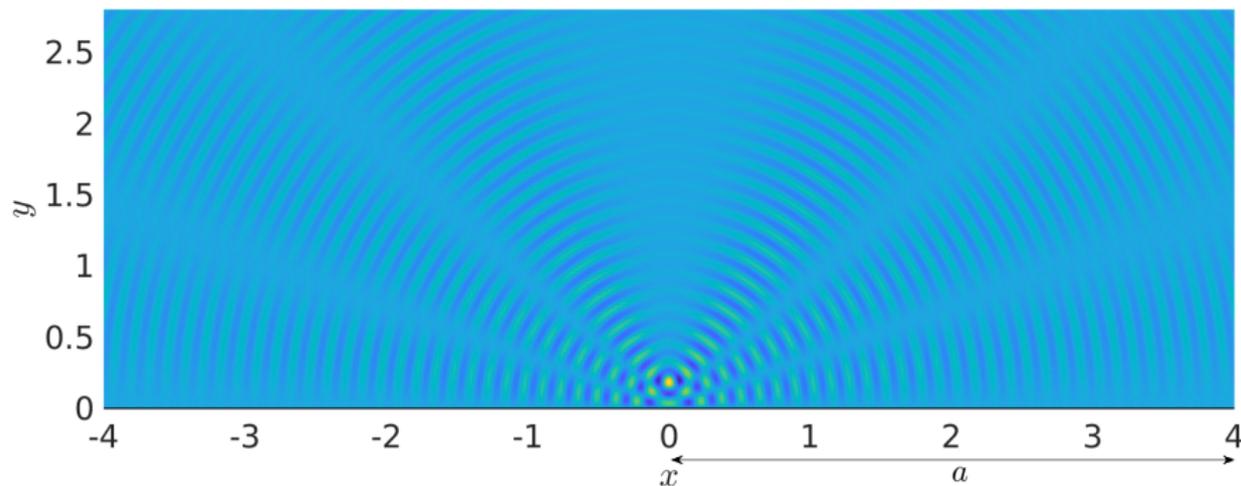
Numerical results: flat  $\Gamma$ :  $f(\tilde{x}) = f_- = 0.25$

2D numerical results when  $\Gamma$  is flat, applying  $h$ -BEM with P1 elements and uniform mesh on part of surface of length  $2a$ , with

$$k = 1, kh = 0.5, kf_- = 0.25, z = (0, 5),$$

using the “Gypsilab” Matlab BEM toolbox of F. Alouges and M. Aussal.

Real part of the total field  $u$

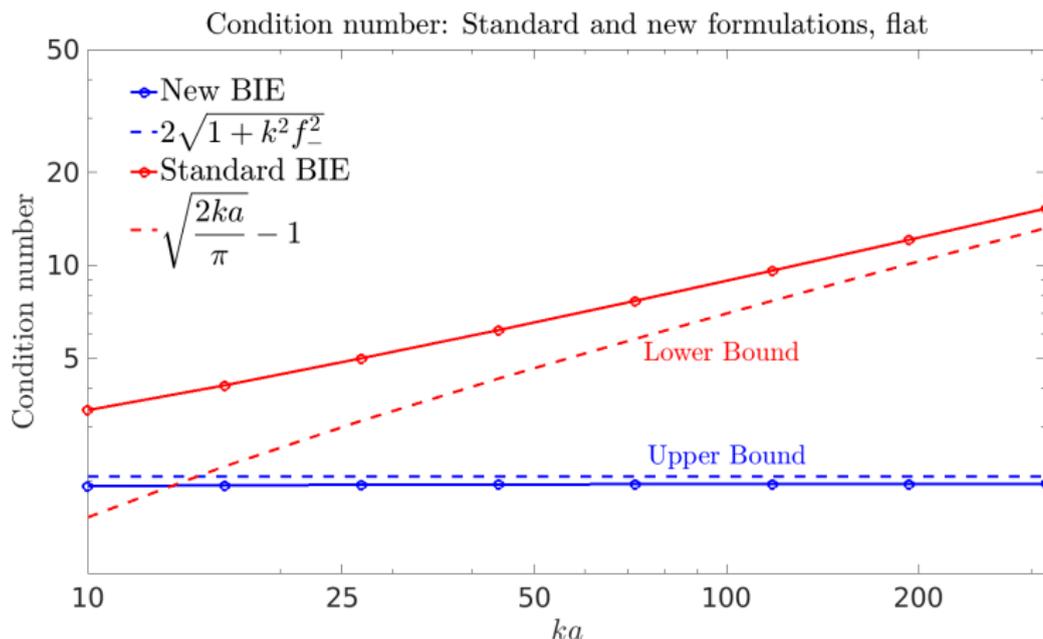


# Numerical results: flat $\Gamma$ : $f(\tilde{x}) = f_- = 0.25$

2D numerical results when  $\Gamma$  is flat, applying  $h$ -BEM with P1 elements and uniform mesh on part of surface of length  $2a$ , with

$$k = 1, kh = 0.5, kf_- = 0.25, z = (0, 5),$$

using the “Gypsilab” Matlab BEM toolbox of F. Alouges and M. Aussal.

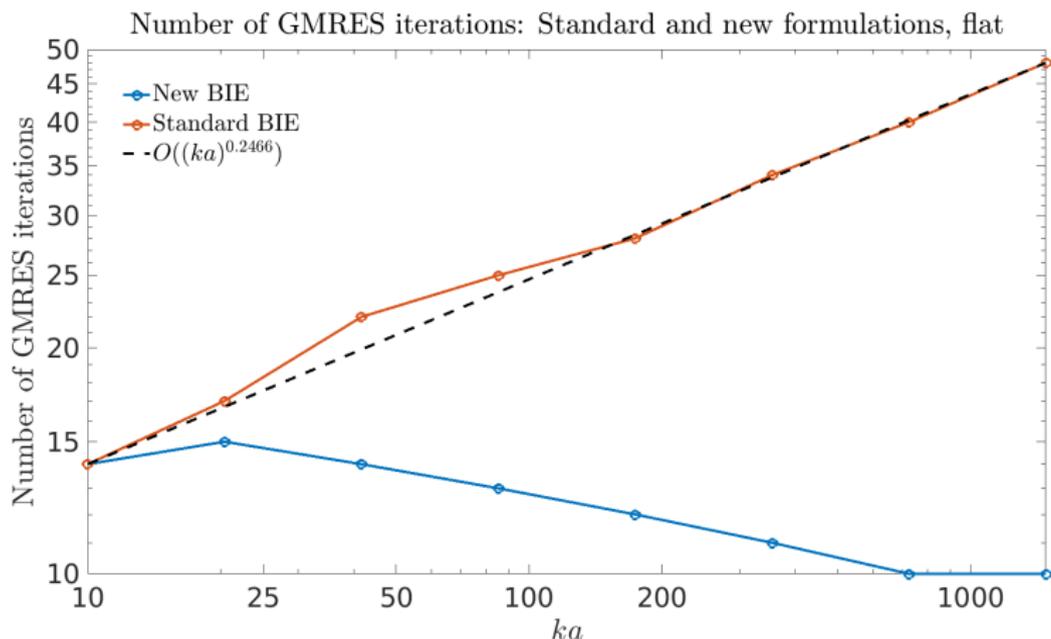


# Numerical results: flat $\Gamma$ : $f(\tilde{x}) = f_- = 0.25$

2D numerical results when  $\Gamma$  is flat, applying  $h$ -BEM with P1 elements and uniform mesh on part of surface of length  $2a$ , with

$$k = 1, kh = 0.5, kf_- = 0.25, z = (0, 5),$$

using the “Gypsilab” Matlab BEM toolbox of F. Alouges and M. Aussal.



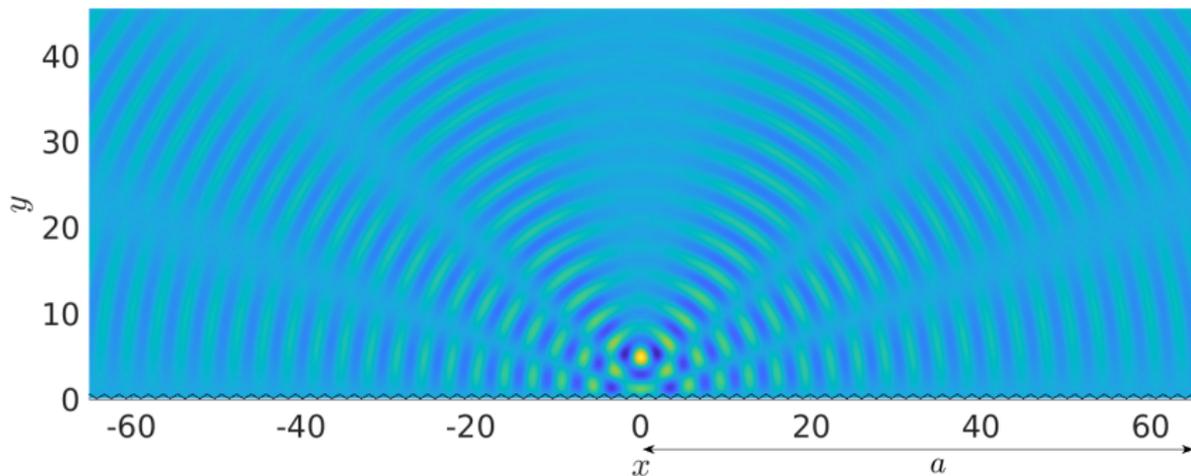
Numerical results: sawtooth  $\Gamma$  :  $f_- \leq f(\tilde{x}) \leq f_+$ , slope  $L$

2D numerical results for sawtooth  $\Gamma$ , applying  $h$ -BEM with P1 elements and uniform mesh on part of surface of length  $2a$ , with

$$k = 2, kh = 0.3, kf_- = 0.25, kf_+ = 1.25 L = 0.578; z = (0, 5),$$

using the “Gypsilab” Matlab BEM toolbox of F. Alouges and M. Aussal.

Real part of the total field  $u$



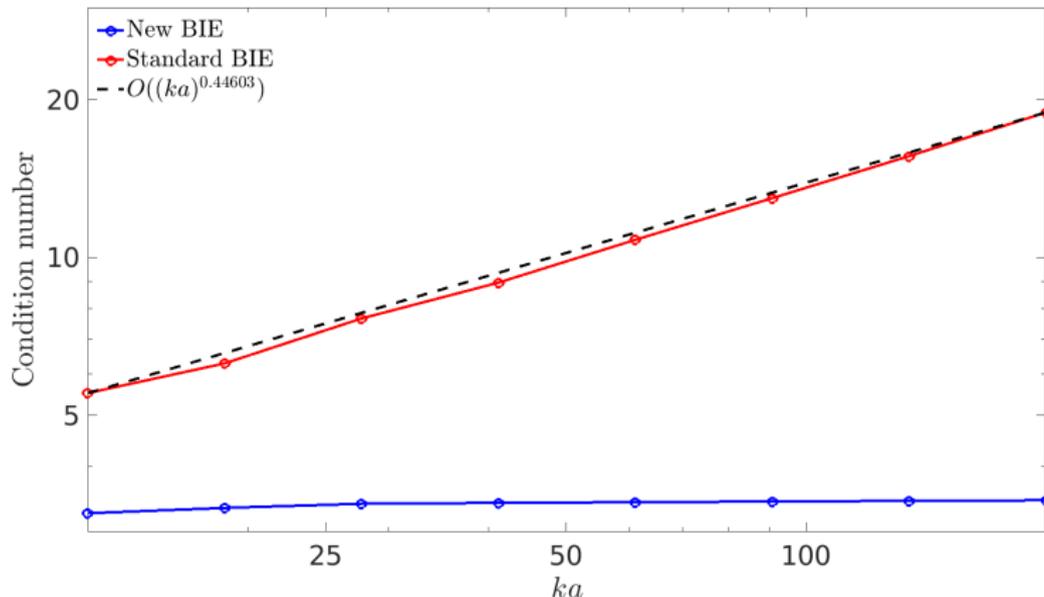
# Numerical results: sawtooth $\Gamma$ : $f_- \leq f(\tilde{x}) \leq f_+$ , slope $L$

2D numerical results for sawtooth  $\Gamma$ , applying  $h$ -BEM with P1 elements and uniform mesh on part of surface of length  $2a$ , with

$$k = 2, kh = 0.3, kf_- = 0.25, kf_+ = 1.25 L = 0.578; z = (0, 5),$$

using the “Gypsilab” Matlab BEM toolbox of F. Alouges and M. Aussal.

Condition number: Standard and new formulations, saw tooth

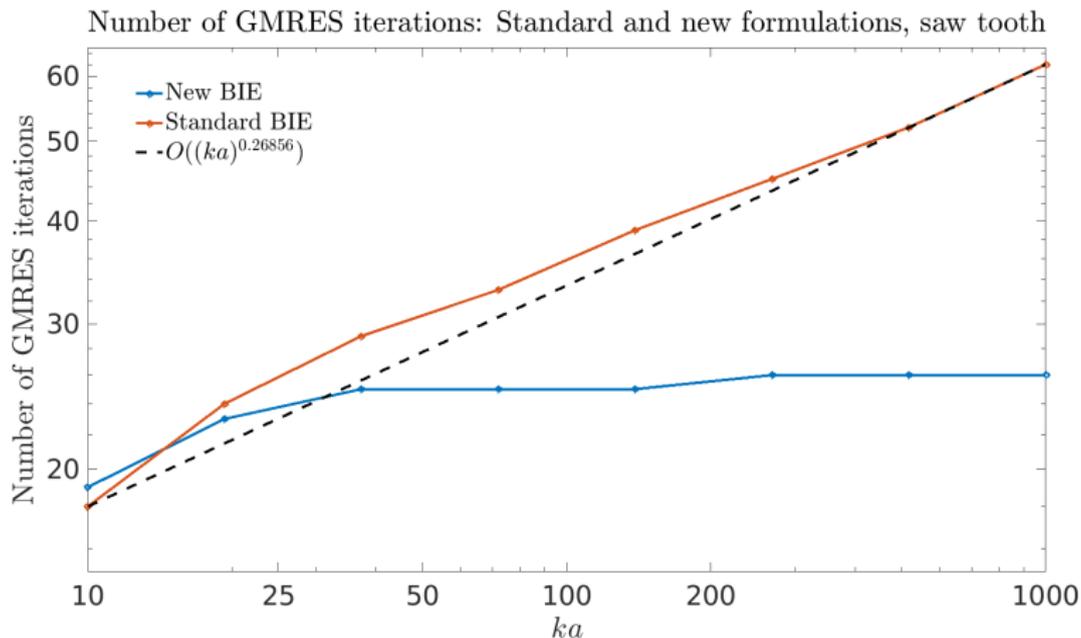


# Numerical results: sawtooth $\Gamma$ : $f_- \leq f(\tilde{x}) \leq f_+$ , slope $L$

2D numerical results for sawtooth  $\Gamma$ , applying  $h$ -BEM with P1 elements and uniform mesh on part of surface of length  $2a$ , with

$$k = 2, kh = 0.3, kf_- = 0.25, kf_+ = 1.25 L = 0.578; z = (0, 5),$$

using the “Gypsilab” Matlab BEM toolbox of F. Alouges and M. Aussal.



# Summary

We have:

- Seen that unbounded rough surfaces are problems with interesting additional computational and numerical analysis challenges!

# Summary

We have:

- Seen that unbounded rough surfaces are problems with interesting additional computational and numerical analysis challenges!
- Recalled the strong/precise results available for analysis of Galerkin methods and GMRES when  $A : H \rightarrow H$  is bounded and coercive

# Summary

We have:

- Seen that unbounded rough surfaces are problems with interesting additional computational and numerical analysis challenges!
- Recalled the strong/precise results available for analysis of Galerkin methods and GMRES when  $A : H \rightarrow H$  is bounded and coercive
- Recalled that, even for bounded obstacles, **no convergence proof exists yet for any Galerkin BEM** for the standard 2nd kind BIE on  $L^2(\Gamma)$  with  $A = \frac{1}{2}I + K' - ikS$ , that applies for general Lipschitz  $\Gamma$ , or even just for all star-shaped polyhedral  $\Gamma$

# Summary

We have:

- Seen that unbounded rough surfaces are problems with interesting additional computational and numerical analysis challenges!
- Recalled the strong/precise results available for analysis of Galerkin methods and GMRES when  $A : H \rightarrow H$  is bounded and coercive
- Recalled that, even for bounded obstacles, **no convergence proof exists yet for any Galerkin BEM** for the standard 2nd kind BIE on  $L^2(\Gamma)$  with  $A = \frac{1}{2}I + K' - ikS$ , that applies for general Lipschitz  $\Gamma$ , or even just for all star-shaped polyhedral  $\Gamma$
- Recalled recent novel 2nd kind integral equations for bounded obstacles, with  $A$  replaced by an operator  $A_Z := Z \cdot n(\frac{1}{2}I + K') + Z \cdot \nabla_{\Gamma} S - ikS$  which is coercive + compact

# Summary

We have:

- Seen that unbounded rough surfaces are problems with interesting additional computational and numerical analysis challenges!
- Recalled the strong/precise results available for analysis of Galerkin methods and GMRES when  $A : H \rightarrow H$  is bounded and coercive
- Recalled that, even for bounded obstacles, **no convergence proof exists yet for any Galerkin BEM** for the standard 2nd kind BIE on  $L^2(\Gamma)$  with  $A = \frac{1}{2}I + K' - ikS$ , that applies for general Lipschitz  $\Gamma$ , or even just for all star-shaped polyhedral  $\Gamma$
- Recalled recent novel 2nd kind integral equations for bounded obstacles, with  $A$  replaced by an operator  $A_Z := Z \cdot n(\frac{1}{2}I + K') + Z \cdot \nabla_{\Gamma} S - ikS$  which is coercive + compact
- Proposed a new 2nd kind integral equation of this type for our RSS problem with  $Z = e_a$ , the constant vertical unit vector, for which  $A_Z$  is **bounded and uniformly-in- $k$  coercive**, leading to proof of **convergence of combined surface truncation/Galerkin BEM**, and **convergence of GMRES** in a number of iterations independent of the element diameter  $h$  and the truncated surface diameter  $a$ .

# References

- D. Baskin, E.A. Spence, J. Wunsch, Sharp high-frequency estimates for the Helmholtz equation and applications to boundary integral equations, *SIAM J. Math. Anal.* 48, 229-267 (2016)
- B. Beckermann, S. A. Goreinov, E. E. Tyrtyshnikov. Some remarks on the Elman estimate for GMRES, *SIAM J. Matrix Anal. Appl.*, 27, 772-778 (2006).
- S. N. Chandler-Wilde, P. Monk, Existence, uniqueness and variational methods for scattering by unbounded rough surfaces, *SIAM J. Math. Anal.* 37, 598-618 (2005)
- S. N. Chandler-Wilde, E. Heinemeyer, R. Potthast, Acoustic scattering by mildly rough unbounded surfaces in three dimensions, *SIAM J. Appl. Math.* 66, 1002-1026 (2006)
- S. N. Chandler-Wilde, E. Heinemeyer, R. Potthast, A well-posed integral equation formulation for 3D rough surface scattering. *Proc. R. Soc. Lond. A* 462, 3683-3705 (2006)
- S. N. Chandler-Wilde, S. Langdon, A Galerkin boundary element method for high frequency scattering by convex polygons, *SIAM J. Numer. Anal.* 43, 610-640 (2007).
- S. N. Chandler-Wilde, I. G. Graham, S. Langdon, M. Lindner, Condition number estimates for combined potential boundary integral operators in acoustic scattering, *J. Integral Equat. Appl.* 21, 229-279 (2009)

- S. N. Chandler-Wilde, J. Elschner, Variational approach in weighted Sobolev spaces to scattering by unbounded rough surfaces, *SIAM J. Math. Anal.* 42, 2554-2580 (2010)
- S. N. Chandler-Wilde, E. A. Spence, Coercivity, essential norms, and the Galerkin method for second-kind integral equations on polyhedral and Lipschitz domains, *Numer. Math.* 150, 299-371 (2022)
- S. N. Chandler-Wilde, E. A. Spence, Coercive second-kind boundary integral equations for the Laplace Dirichlet problem on Lipschitz domains, *arXiv:2210.02432*, 2022.
- A. Gibbs, S. N. Chandler-Wilde, S. Langdon, A. Moiola, A high frequency boundary element method for scattering by a class of multiple obstacles, *IMA J. Numer. Anal.* 41, 1197-1239 (2021)
- V. Y. Gotlib, Solutions of the Helmholtz equation, concentrated near a plane periodic boundary, *J. Math. Sci.* 102, 4188-4194 (2000)
- A. Rathsfeld, Simulating rough surfaces by periodic and biperiodic gratings, *Weierstrass Institute Preprint*, Berlin, 2022
- E. A. Spence, S. N. Chandler-Wilde, I. G. Graham, V. P. Smyshlyaev, A new frequency-uniform coercive boundary integral equation for acoustic scattering, *Comm. Pure Appl. Math.* 64, 1384-1415 (2011)
- E. A. Spence, I. V. Kamotski, V. P. Smyshlyaev, Coercivity of combined boundary integral equations in high-frequency scattering, *Comm. Pure Appl. Math.* 68, 1587-1639 (2015)

- G. Verchota, Layer potentials and regularity for the Dirichlet problem for Laplace's equation in Lipschitz domains, *J. Funct. Anal.* 59, 572-611 (1984)
- B. Zhang, S. N. Chandler-Wilde, Integral equation methods for scattering by infinite rough surfaces. *Math. Methods Appl. Sci.* 26, 463-488 (2003)