

Wave scattering by trapping obstacles: resolvent estimates and applications to boundary integral equations and their numerical solution

Simon Chandler-Wilde

Department of Mathematics and
Statistics
University of Reading
s.n.chandler-wilde@reading.ac.uk



Joint work with:

Euan Spence (Bath), Andrew Gibbs (Reading/Leuven), Valery Smyshlyaev (UCL)

Analysis (Applied and PDE) Seminar:
Heriot-Watt University, September 2017

More info: [new preprint "High-frequency bounds ..." on arXiv](#)

The wave equation and Helmholtz equation

In acoustics the increase in air pressure at x at time t , $U(x, t)$, satisfies

$$\Delta U = \frac{1}{c^2} \frac{\partial^2 U}{\partial t^2} \quad \left(\Delta = \frac{\partial^2}{\partial x_1^2} + \frac{\partial^2}{\partial x_2^2} + \frac{\partial^2}{\partial x_3^2} \right).$$

The wave equation and Helmholtz equation

In acoustics the increase in air pressure at x at time t , $U(x, t)$, satisfies

$$\Delta U = \frac{1}{c^2} \frac{\partial^2 U}{\partial t^2} \quad \left(\Delta = \frac{\partial^2}{\partial x_1^2} + \frac{\partial^2}{\partial x_2^2} + \frac{\partial^2}{\partial x_3^2} \right).$$

If time-dependence is **time harmonic**, i.e.,

$$U(x, t) = A(x) \cos(\phi(x) - \omega t),$$

for some $\omega = 2\pi f > 0$, with $f =$ **frequency**

The wave equation and Helmholtz equation

In acoustics the increase in air pressure at x at time t , $U(x, t)$, satisfies

$$\Delta U = \frac{1}{c^2} \frac{\partial^2 U}{\partial t^2} \quad \left(\Delta = \frac{\partial^2}{\partial x_1^2} + \frac{\partial^2}{\partial x_2^2} + \frac{\partial^2}{\partial x_3^2} \right).$$

If time-dependence is **time harmonic**, i.e.,

$$U(x, t) = A(x) \cos(\phi(x) - \omega t),$$

for some $\omega = 2\pi f > 0$, with $f = \mathbf{frequency}$, then

$$U(x, t) = \Re(u(x)e^{-i\omega t})$$

where $u(x) = A(x) \exp(i\phi(x))$ satisfies the **Helmholtz equation**

$$\Delta u + k^2 u = 0,$$

with $k = \omega/c$ the **wavenumber**.

The wave equation and Helmholtz equation

In acoustics the increase in air pressure at x at time t , $U(x, t)$, satisfies

$$\Delta U = \frac{1}{c^2} \frac{\partial^2 U}{\partial t^2} \quad \left(\Delta = \frac{\partial^2}{\partial x_1^2} + \frac{\partial^2}{\partial x_2^2} + \frac{\partial^2}{\partial x_3^2} \right).$$

If time-dependence is **time harmonic**, i.e.,

$$U(x, t) = A(x) \cos(\phi(x) - \omega t),$$

for some $\omega = 2\pi f > 0$, with $f = \mathbf{frequency}$, then

$$U(x, t) = \Re(u(x)e^{-i\omega t})$$

where $u(x) = A(x) \exp(i\phi(x))$ satisfies the **Helmholtz equation**

$$\Delta u + k^2 u = 0,$$

with $k = \omega/c$ the **wavenumber**. E.g. if $u(x) = \exp(ikx \cdot d)$, for some **unit vector** d , then

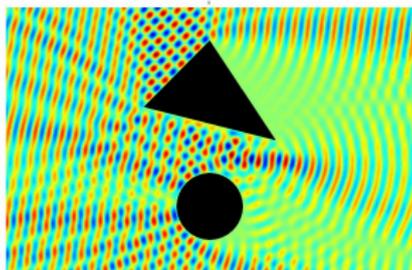
$$U(x, t) = \Re(u(x)e^{-i\omega t}) = \cos(kx \cdot d - \omega t)$$

is a **plane wave** travelling in direction d with **wavelength**

$$\lambda = 2\pi/k = c/f.$$

Challenges of $\Delta u + k^2 u = 0$ when k is large

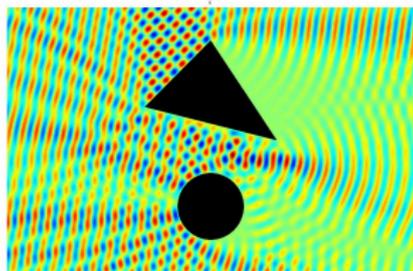
1. Solution is oscillatory and multiscale: one scale is the wavelength $\lambda = 2\pi/k$.



$\Re(u(x)) = U(x, 0)$ for 2-d scattering of incident plane wave
 $u^{\text{inc}}(x) = \exp(ikd \cdot x)$.

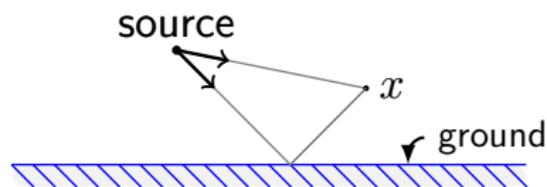
Challenges of $\Delta u + k^2 u = 0$ when k is large

1. Solution is oscillatory and multiscale: one scale is the wavelength $\lambda = 2\pi/k$.



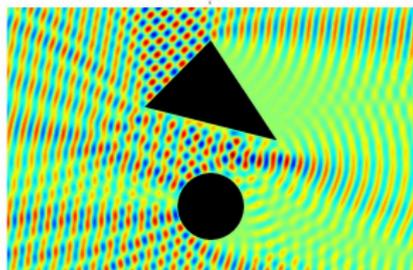
$\Re(u(x)) = U(x, 0)$ for 2-d scattering of incident plane wave $u^{\text{inc}}(x) = \exp(ikd \cdot x)$.

2. In the **singular limit** $k \rightarrow \infty$ the wave equation transitions to a particle/ray/billiards model



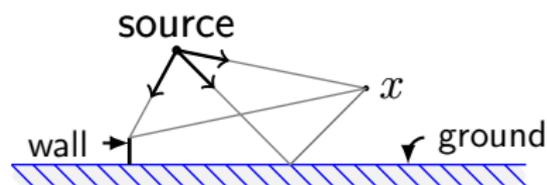
Challenges of $\Delta u + k^2 u = 0$ when k is large

1. Solution is oscillatory and multiscale: one scale is the wavelength $\lambda = 2\pi/k$.



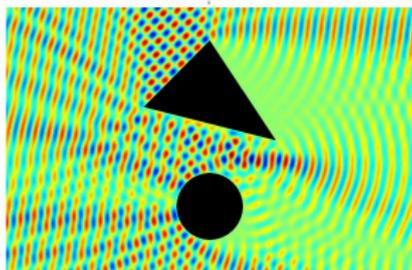
$\Re(u(x)) = U(x, 0)$ for 2-d scattering of incident plane wave $u^{\text{inc}}(x) = \exp(ikd \cdot x)$.

2. In the **singular limit** $k \rightarrow \infty$ the wave equation transitions to a particle/ray/billiards model



Challenges of $\Delta u + k^2 u = 0$ when k is large

1. Solution is oscillatory and multiscale: one scale is the wavelength $\lambda = 2\pi/k$.



$\Re(u(x)) = U(x, 0)$ for 2-d scattering of incident plane wave $u^{\text{inc}}(x) = \exp(ikd \cdot x)$.

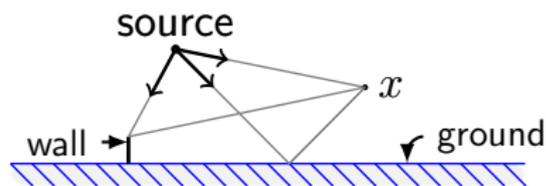
2. In the **singular limit** $k \rightarrow \infty$ the wave equation transitions to a particle/ray/billiards model

$$u(x) \approx \sum_j u_j(x)$$

where sum over **rays** passing through x , with

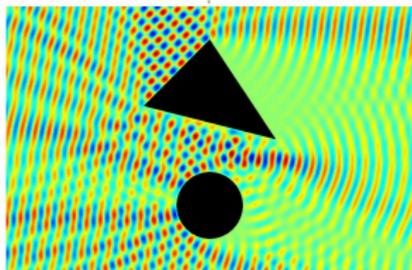
$\arg u_j(x) =$ **optical length** of ray path $= ks_j$

$|u_j(x)| =$ **amplitude** determined by energy conservation



Challenges of $\Delta u + k^2 u = 0$ when k is large

1. Solution is oscillatory and multiscale: one scale is the wavelength $\lambda = 2\pi/k$.



$\Re(u(x)) = U(x, 0)$ for 2-d scattering of incident plane wave $u^{\text{inc}}(x) = \exp(ikd \cdot x)$.

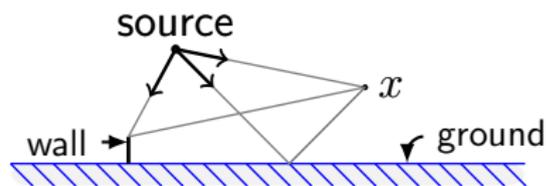
2. In the **singular limit** $k \rightarrow \infty$ the wave equation transitions to a particle/ray/billiards model

$$u(x) \approx \sum_j u_j(x)$$

where sum over **rays** passing through x , with

$\arg u_j(x) =$ **optical length** of ray path $= k s_j$

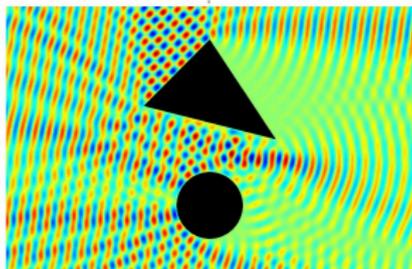
$|u_j(x)| =$ **amplitude** determined by energy conservation



but with multiplication of $u_j(x)$ by **coefficients** accounting for **reflection** and **diffraction** events.

Challenges of $\Delta u + k^2 u = 0$ when k is large

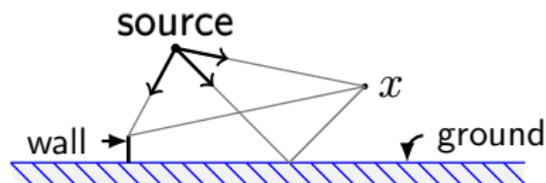
1. Solution is oscillatory and multiscale: one scale is the wavelength $\lambda = 2\pi/k$.



$\Re(u(x)) = U(x, 0)$ for 2-d scattering of incident plane wave $u^{\text{inc}}(x) = \exp(ikd \cdot x)$.

2. In the **singular limit** $k \rightarrow \infty$ the wave equation transitions to a particle/ray/billiards model

$$u(x) \approx \sum_j u_j(x)$$



where sum over **rays** passing through x , with

$\arg u_j(x) =$ **optical length** of ray path $= ks_j$

$|u_j(x)| =$ **amplitude** determined by energy conservation

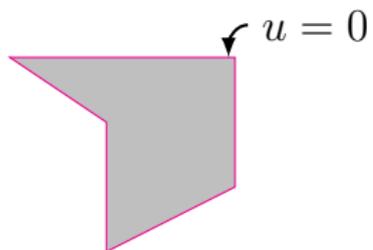
The rigorous justification of such approximations is the concern of **semi-classical** analysis.

but with multiplication of $u_j(x)$ by **coefficients** accounting for **reflection** and **diffraction** events.

What is this talk about?

u satisfies Sommerfeld rad. cond. (SRC)

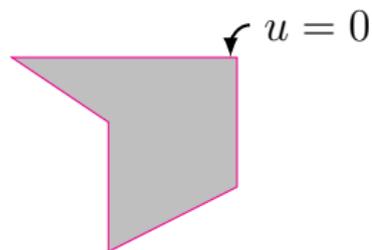
$$\Delta u + k^2 u = f \text{ (source, compactly supported)}$$



What is this talk about?

u satisfies Sommerfeld rad. cond. (SRC)

$$\Delta u + k^2 u = f \text{ (source, compactly supported)}$$

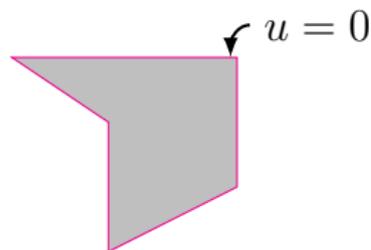


This talk is about **wavenumber-explicit** bounds, i.e. bounds explicit in k , for wave scattering obstacles: focus on sound soft (Dirichlet) case and **large** k .

What is this talk about?

u satisfies Sommerfeld rad. cond. (SRC)

$$\Delta u + k^2 u = f \text{ (source, compactly supported)}$$



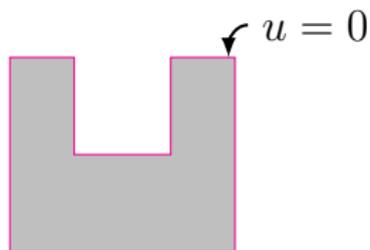
This talk is about **wavenumber-explicit** bounds, i.e. bounds explicit in k , for wave scattering obstacles: focus on sound soft (Dirichlet) case and **large** k .

It's about cases where the obstacle is **nontrapping**, e.g. **star-shaped** (like above example).

What is this talk about?

u satisfies Sommerfeld rad. cond. (SRC)

$$\Delta u + k^2 u = f \text{ (source, compactly supported)}$$



This talk is about **wavenumber-explicit** bounds, i.e. bounds explicit in k , for wave scattering obstacles: focus on sound soft (Dirichlet) case and **large** k .

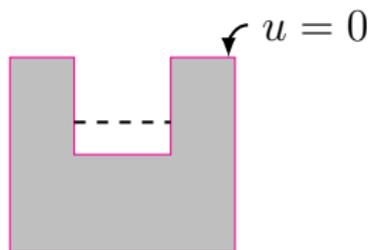
It's about cases where the obstacle is **nontrapping**, e.g. **star-shaped** (like above example).

But particularly about cases where the obstacle is **trapping**

What is this talk about?

u satisfies Sommerfeld rad. cond. (SRC)

$$\Delta u + k^2 u = f \text{ (source, compactly supported)}$$



This talk is about **wavenumber-explicit** bounds, i.e. bounds explicit in k , for wave scattering obstacles: focus on sound soft (Dirichlet) case and **large** k .

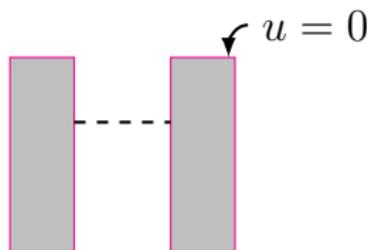
It's about cases where the obstacle is **nontrapping**, e.g. **star-shaped** (like above example).

But particularly about cases where the obstacle is **trapping** supporting a **trapped ray/billiard trajectory**.

What is this talk about?

u satisfies Sommerfeld rad. cond. (SRC)

$$\Delta u + k^2 u = f \text{ (source, compactly supported)}$$



This talk is about **wavenumber-explicit** bounds, i.e. bounds explicit in k , for wave scattering obstacles: focus on sound soft (Dirichlet) case and **large k** .

It's about cases where the obstacle is **nontrapping**, e.g. **star-shaped** (like above example).

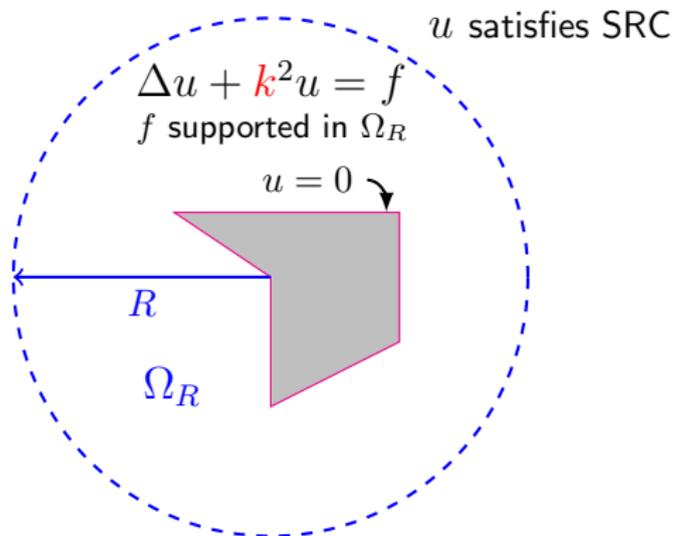
But particularly about cases where the obstacle is **trapping** supporting a **trapped ray/billiard trajectory**.

Including cases where the obstacle has more than one component, in other words **multiple scattering**.

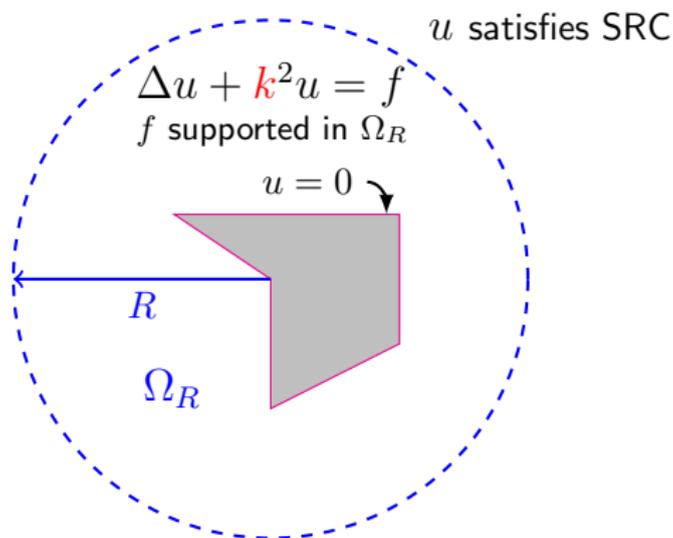
Overview of Talk

- 1 Helmholtz equation: what is it and why interesting
- 2 What is this talk about?
- 3 Resolvent estimates
 - What are they?
 - The three known estimates and their geometries
 - A new estimate for parabolic trapping
 - The Morawetz/Rellich identity method of proof
- 4 Implications for Boundary Integral Equations
- 5 Implications for hp -BEM
- 6 Conclusions

What is a resolvent estimate?



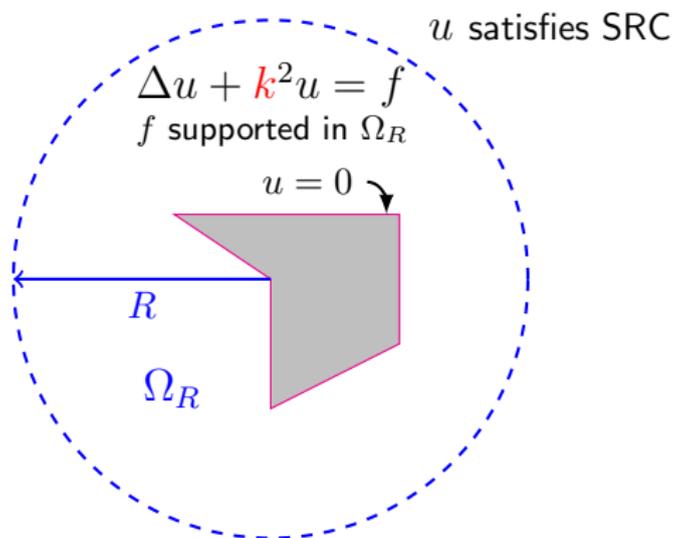
What is a resolvent estimate?



It is the wavenumber-explicit bound that, for $R > 0$, and some specified $c(k)$,

$$\|\nabla u\|_{L^2(\Omega_R)} + k\|u\|_{L^2(\Omega_R)} \lesssim c(k)\|f\|_{L^2(\Omega_R)}, \quad \text{for } k \geq k_0 > 0.$$

What is a resolvent estimate?

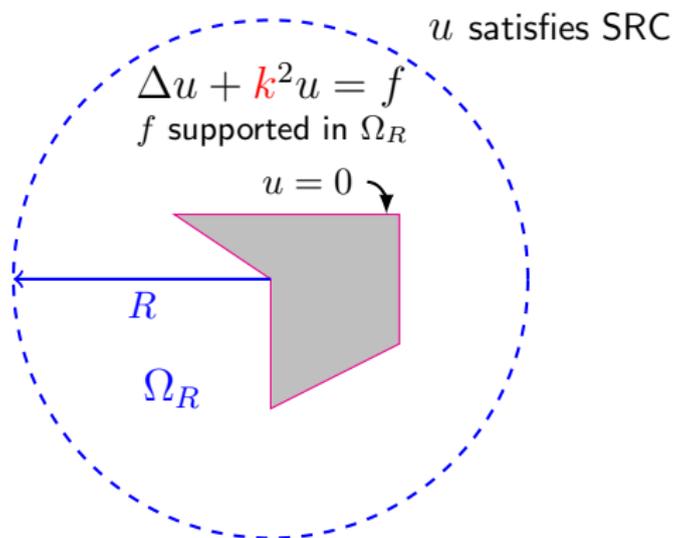


It is the wavenumber-explicit bound that, for $R > 0$, and some specified $c(k)$,

$$\|\nabla u\|_{L^2(\Omega_R)} + k\|u\|_{L^2(\Omega_R)} \lesssim c(k)\|f\|_{L^2(\Omega_R)}, \quad \text{for } k \geq k_0 > 0.$$

$A \lesssim B$ means $A \leq CB$, where $C > 0$ independent of k and f , but depends on R .

What is a resolvent estimate?



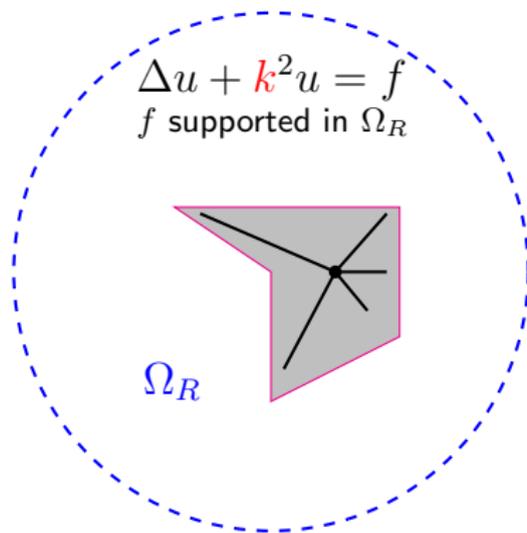
It is the wavenumber-explicit bound that, for $R > 0$, and some specified $c(k)$,

$$\|\nabla u\|_{L^2(\Omega_R)} + k\|u\|_{L^2(\Omega_R)} \lesssim c(k)\|f\|_{L^2(\Omega_R)}, \quad \text{for } k \geq k_0 > 0.$$

$A \lesssim B$ means $A \leq CB$, where $C > 0$ independent of k and f , but depends on R .

We will see that resolvent estimates give us: bounds on **DtN maps**, on inverses of **boundary integral operators**, on errors in **FEM**, **BEM**, ...

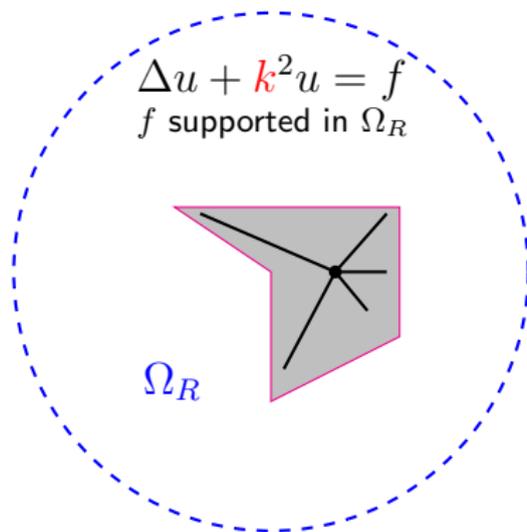
The known estimates and their geometries



Star-shaped obstacle (C^∞ : Morawetz 1975; C^0 : C-W & Monk 2008)

$$\|\nabla u\|_{L^2(\Omega_R)} + k \|u\|_{L^2(\Omega_R)} \lesssim \|f\|_{L^2(\Omega_R)}, \quad \text{i.e. } c(k) = 1$$

The known estimates and their geometries

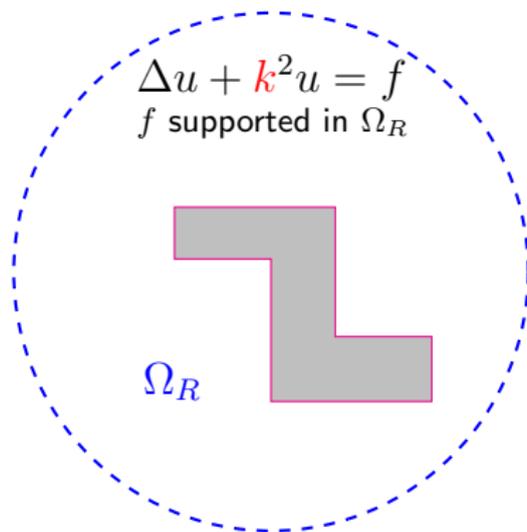


Star-shaped obstacle (C^∞ : Morawetz 1975; C^0 : C-W & Monk 2008)

$$\|\nabla u\|_{L^2(\Omega_R)} + k\|u\|_{L^2(\Omega_R)} \lesssim \|f\|_{L^2(\Omega_R)}, \quad \text{i.e. } c(k) = 1$$

Best possible bound: achieved by $u(x) = \chi(x) \exp(ikx_1)$, if $\chi \in C_0^\infty(\Omega_R)$.

The known estimates and their geometries

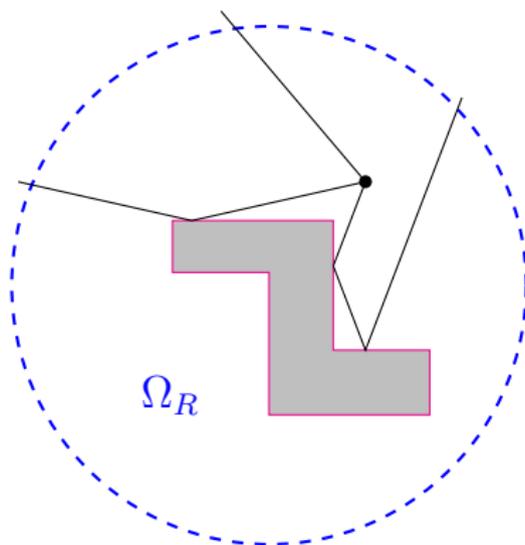


Nontrapping obstacle (C^∞ : Morawetz, Ralston, Strauss 1977, Vainberg 1975, Melrose & Sjöstrand 1982; polygon: Baskin & Wunsch 2013)

$$\|\nabla u\|_{L^2(\Omega_R)} + k\|u\|_{L^2(\Omega_R)} \lesssim \|f\|_{L^2(\Omega_R)}, \quad \text{i.e. } c(k) = 1$$

Best possible bound: achieved by $u(x) = \chi(x) \exp(ikx_1)$, if $\chi \in C_0^\infty(\Omega_R)$.

The known estimates and their geometries

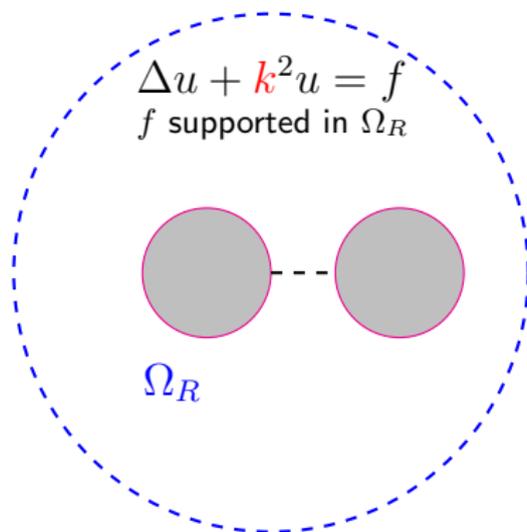


Nontrapping obstacle (C^∞ : Morawetz, Ralston, Strauss 1977, Vainberg 1975, Melrose & Sjöstrand 1982; polygon: Baskin & Wunsch 2013)

$$\|\nabla u\|_{L^2(\Omega_R)} + k\|u\|_{L^2(\Omega_R)} \lesssim \|f\|_{L^2(\Omega_R)}, \quad \text{i.e. } c(k) = 1$$

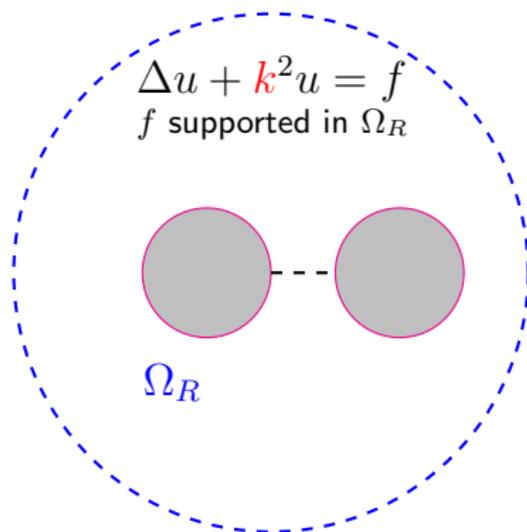
Nontrapping: there exists $T > 0$ such that all the billiard trajectories starting in Ω_R at time zero and travelling at unit speed leave Ω_R by time T .

The known estimates and their geometries



Two or more C^∞ strictly convex, positive curvature obstacles (Ikawa 1988, Burq 2004), example of **hyperbolic, unstable trapping**

The known estimates and their geometries

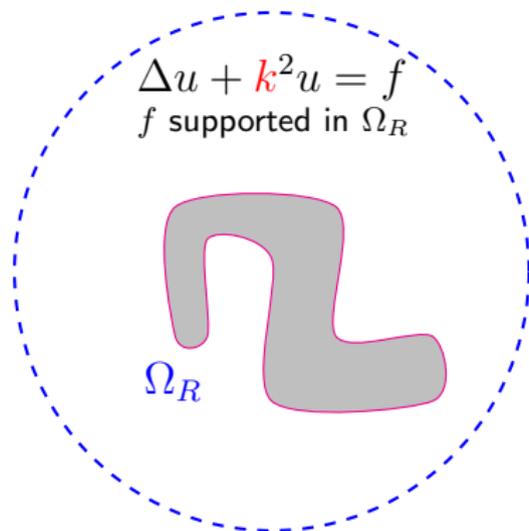


Two or more C^∞ strictly convex, positive curvature obstacles (Ikawa 1988, Burq 2004), example of **hyperbolic, unstable trapping**

$$\|\nabla u\|_{L^2(\Omega_R)} + k\|u\|_{L^2(\Omega_R)} \lesssim \log(2+k)\|f\|_{L^2(\Omega_R)}, \quad \text{i.e. } c(k) = \log(2+k),$$

so **only logarithmically worse** than the nontrapping case.

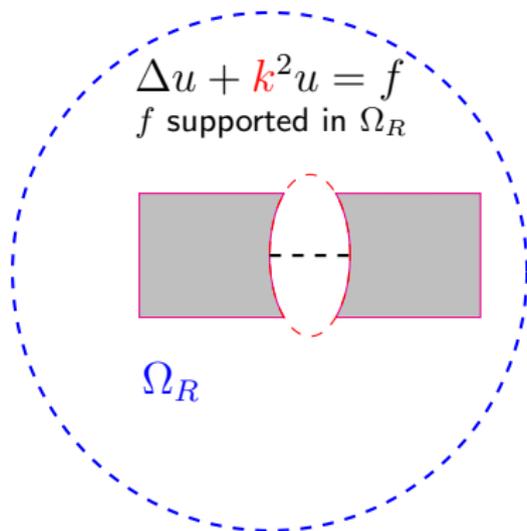
The known estimates and their geometries



General C^∞ “**worst case**” bound (Burq 1998): for some $\alpha > 0$,

$$\|\nabla u\|_{L^2(\Omega_R)} + k\|u\|_{L^2(\Omega_R)} \lesssim \exp(\alpha k)\|f\|_{L^2(\Omega_R)}, \quad \text{i.e. } c(k) = \exp(\alpha k).$$

The known estimates and their geometries

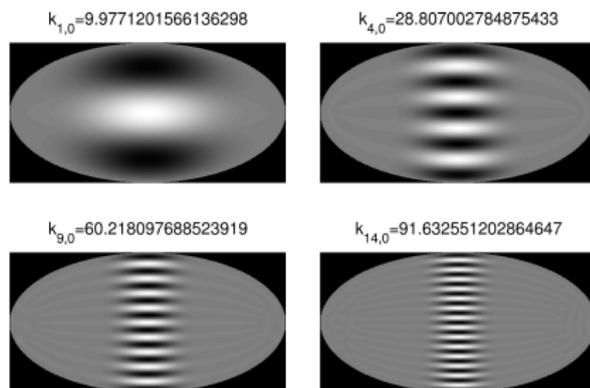
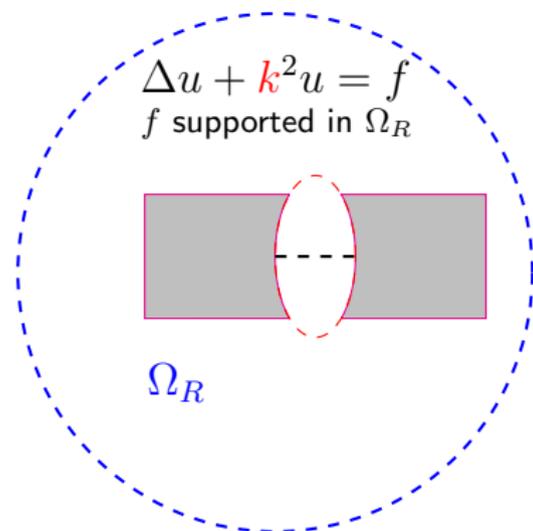


General C^∞ “**worst case**” bound (Burq 1998): for some $\alpha > 0$,

$$\|\nabla u\|_{L^2(\Omega_R)} + k\|u\|_{L^2(\Omega_R)} \lesssim \exp(\alpha k)\|f\|_{L^2(\Omega_R)}, \quad \text{i.e. } c(k) = \exp(\alpha k).$$

This achieved for some $k_m \rightarrow \infty$ when there is **elliptic, stable trapping** (Cardoso, Popov 2002; Betcke, C-W, Graham, Langdon, Lindner 2011) with a **quasimode localised around the trapped ray**.

The known estimates and their geometries



General C^∞ “**worst case**” bound (Burq 1998): for some $\alpha > 0$,

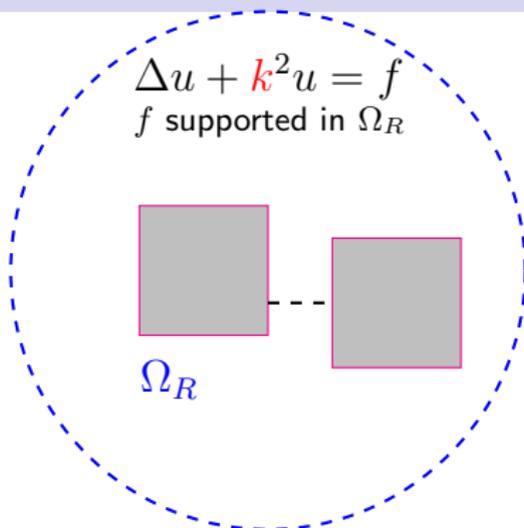
$$\|\nabla u\|_{L^2(\Omega_R)} + k\|u\|_{L^2(\Omega_R)} \lesssim \exp(\alpha k)\|f\|_{L^2(\Omega_R)}, \quad \text{i.e. } c(k) = \exp(\alpha k).$$

This achieved for some $k_m \rightarrow \infty$ when there is **elliptic, stable trapping** (Cardoso, Popov 2002; Betcke, C-W, Graham, Langdon, Lindner 2011) with a **quasimode localised around the trapped ray**.

Where have we got to in the talk?

- 1 Helmholtz equation: what is it and why interesting
- 2 What is this talk about?
- 3 Resolvent estimates
 - What are they?
 - The three known estimates and their geometries
 - **A new estimate for parabolic trapping**
 - The Morawetz/Rellich identity method of proof
- 4 Implications for Boundary Integral Equations
- 5 Implications for hp -BEM
- 6 Conclusions

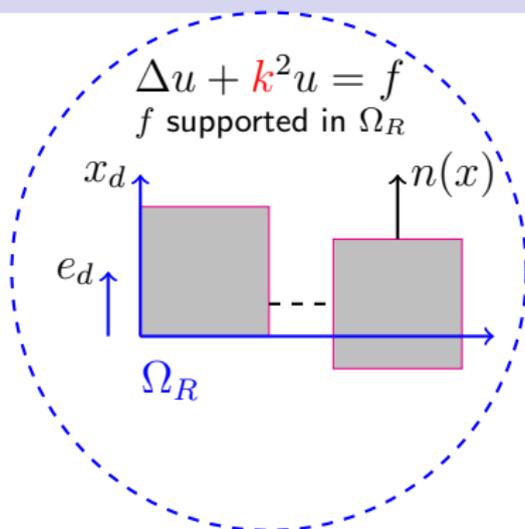
Our new estimate for parabolic, neutral trapping



Theorem (C-W, Spence, Gibbs, Smyshlyaev 2017)

$$\|\nabla u\|_{L^2(\Omega_R)} + k\|u\|_{L^2(\Omega_R)} \lesssim k^2\|f\|_{L^2(\Omega_R)}, \quad \text{i.e. } c(k) = k^2.$$

Our new estimate for parabolic, neutral trapping



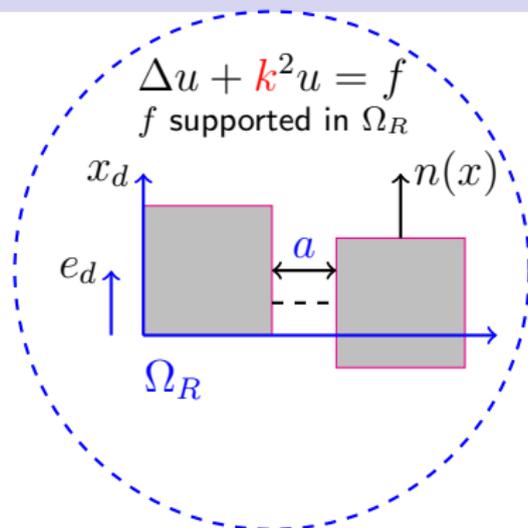
Theorem (C-W, Spence, Gibbs, Smyshlyaev 2017)

$$\|\nabla u\|_{L^2(\Omega_R)} + k\|u\|_{L^2(\Omega_R)} \lesssim k^2\|f\|_{L^2(\Omega_R)}, \quad \text{i.e. } c(k) = k^2.$$

Applies to a general Lipschitz obstacle class, in particular when

$$x_d e_d \cdot n(x) \geq 0 \quad \text{on the boundary}$$

Our new estimate for parabolic, neutral trapping



Theorem (C-W, Spence, Gibbs, Smyshlyaev 2017)

$$\|\nabla u\|_{L^2(\Omega_R)} + k\|u\|_{L^2(\Omega_R)} \lesssim k^2\|f\|_{L^2(\Omega_R)}, \quad \text{i.e. } c(k) = k^2.$$

Applies to a general Lipschitz obstacle class, in particular when

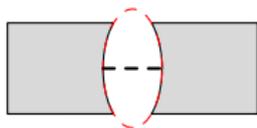
$$x_d e_d \cdot n(x) \geq 0 \quad \text{on the boundary}$$

Further, $\|\nabla u\|_{L^2(\Omega_R)} + k\|u\|_{L^2(\Omega_R)} \gtrsim k\|f\|_{L^2(\Omega_R)}$, for $k = m\pi/a$, $m = 1, 2, \dots$

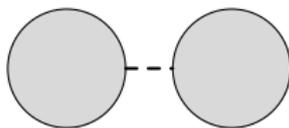
Recap of resolvent estimates for trapping obstacles

$$\|\nabla u\|_{L^2(\Omega_R)} + k\|u\|_{L^2(\Omega_R)} \lesssim c(k)\|f\|_{L^2(\Omega_R)}, \quad \text{for } k \geq k_0 > 0,$$

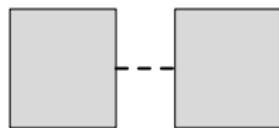
where $c(k) = 1$ for **nontrapping** obstacles, and



$c(k) = \exp(\alpha k)$
elliptic



$c(k) = \log(2 + k)$
hyperbolic



$c(k) = k^2$
parabolic

How are these resolvent estimates proved?

How are these resolvent estimates proved?

The Morawetz/Rellich identity method

Used for:

- Star-shaped obstacles (Morawetz 1975, C-W, Monk 2008)
- “Nearly all” nontrapping obstacles in 2-d (Morawetz, Ralston, Strauss 1977)
- A class of parabolic trapping obstacles (C-W, Spence, Gibbs, Smyshlyaev 2017)

How are these resolvent estimates proved?

The Morawetz/Rellich identity method

Used for:

- Star-shaped obstacles (Morawetz 1975, C-W, Monk 2008)
- “Nearly all” nontrapping obstacles in 2-d (Morawetz, Ralston, Strauss 1977)
- A class of parabolic trapping obstacles (C-W, Spence, Gibbs, Smyshlyaev 2017)

Cathleen Morawetz (1923-2017), Courant Institute, New York.



Listen to the interviews at

<https://www.simonsfoundation.org/2012/12/20/cathleen-morawetz/>

e.g. on women in mathematics, working with Courant, Courant and flexible working, the founding of the Courant Institute, ...

How are these resolvent estimates proved?

The Morawetz/Rellich identity method

Used for:

- Star-shaped obstacles (Morawetz 1975, C-W, Monk 2008)
- “Nearly all” nontrapping obstacles in 2-d (Morawetz, Ralston, Strauss 1977)
- A class of parabolic trapping obstacles (C-W, Spence, Gibbs, Smyshlyaev 2017)

How are these resolvent estimates proved?

The Morawetz/Rellich identity method

Used for:

- Star-shaped obstacles (Morawetz 1975, C-W, Monk 2008)
- “Nearly all” nontrapping obstacles in 2-d (Morawetz, Ralston, Strauss 1977)
- A class of parabolic trapping obstacles (C-W, Spence, Gibbs, Smyshlyaev 2017)

Define **Morawetz multiplier** $\mathcal{Z}u$ by

$$\mathcal{Z}u := Z \cdot \nabla u - ik\beta u + \alpha u,$$

where Z , α , β are real-valued, with $Z \cdot n \geq 0$ on the boundary

How are these resolvent estimates proved?

The Morawetz/Rellich identity method

Used for:

- Star-shaped obstacles (Morawetz 1975, C-W, Monk 2008)
- “Nearly all” nontrapping obstacles in 2-d (Morawetz, Ralston, Strauss 1977)
- A class of parabolic trapping obstacles (C-W, Spence, Gibbs, Smyshlyaev 2017)

Define **Morawetz multiplier** $\mathcal{Z}u$ by

$$\mathcal{Z}u := Z \cdot \nabla u - ik\beta u + \alpha u,$$

where Z , α , β are real-valued, with $Z \cdot n \geq 0$ on the boundary, and

$$2\Re \int_{\Omega_R} \overline{\mathcal{Z}u} f \, dx = 2\Re \int_{\Omega_R} \overline{\mathcal{Z}u} (\Delta u + k^2 u) \, dx = \int_{\Omega_R} +ve + \int_{\partial\Omega_R} +ve$$

How are these resolvent estimates proved?

The Morawetz/Rellich identity method

Used for:

- Star-shaped obstacles (Morawetz 1975, C-W, Monk 2008)
- “Nearly all” nontrapping obstacles in 2-d (Morawetz, Ralston, Strauss 1977)
- A class of parabolic trapping obstacles (C-W, Spence, Gibbs, Smyshlyaev 2017)

Define **Morawetz multiplier** $\mathcal{Z}u$ by

$$\mathcal{Z}u := Z \cdot \nabla u - ik\beta u + \alpha u,$$

where Z , α , β are real-valued, with $Z \cdot n \geq 0$ on the boundary, and

$$2\Re \int_{\Omega_R} \overline{\mathcal{Z}u} f \, dx = 2\Re \int_{\Omega_R} \overline{\mathcal{Z}u} (\Delta u + k^2 u) \, dx = \int_{\Omega_R} +ve + \int_{\partial\Omega_R} +ve$$

For **star-shaped** obstacles use $Z(x) = x$, $\alpha = (d-1)/2$, and $\beta(x) = |x|$ (Morawetz) or $\beta = R$ (C-W/Monk), to get

$$\int_{\Omega_R} (|\nabla u|^2 + k^2 |u|^2) \, dx = -2\Re \int_{\Omega_R} \overline{\mathcal{Z}u} f \, dx - \int_{\partial\Omega_R} +ve \leq \epsilon \|\mathcal{Z}u\|_{L^2(\Omega_R)}^2 + \epsilon^{-1} \|f\|_{L^2(\Omega_R)}^2.$$

How are these resolvent estimates proved?

The Morawetz/Rellich identity method

Used for:

- Star-shaped obstacles (Morawetz 1975, C-W, Monk 2008)
- “Nearly all” nontrapping obstacles in 2-d (Morawetz, Ralston, Strauss 1977)
- A class of parabolic trapping obstacles (C-W, Spence, Gibbs, Smyshlyaev 2017)

Define **Morawetz multiplier** $\mathcal{Z}u$ by

$$\mathcal{Z}u := Z \cdot \nabla u - ik\beta u + \alpha u,$$

where Z , α , β are real-valued, with $Z \cdot n \geq 0$ on the boundary, and

$$2\Re \int_{\Omega_R} \overline{\mathcal{Z}u} f \, dx = 2\Re \int_{\Omega_R} \overline{\mathcal{Z}u} (\Delta u + k^2 u) \, dx = \int_{\Omega_R} +ve + \int_{\partial\Omega_R} +ve$$

For **star-shaped** obstacles use $Z(x) = x$, $\alpha = (d-1)/2$, and $\beta(x) = |x|$ (Morawetz) or $\beta = R$ (C-W/Monk), to get

$$\int_{\Omega_R} (|\nabla u|^2 + k^2 |u|^2) \, dx = -2\Re \int_{\Omega_R} \overline{\mathcal{Z}u} f \, dx - \int_{\partial\Omega_R} +ve \leq \epsilon \|\mathcal{Z}u\|_{L^2(\Omega_R)}^2 + \epsilon^{-1} \|f\|_{L^2(\Omega_R)}^2.$$

How are these resolvent estimates proved?

The Morawetz/Rellich identity method

Used for:

- Star-shaped obstacles (Morawetz 1975, C-W, Monk 2008)
- “Nearly all” nontrapping obstacles in 2-d (Morawetz, Ralston, Strauss 1977)
- A class of parabolic trapping obstacles (C-W, Spence, Gibbs, Smyshlyaev 2017)

Define **Morawetz multiplier** $\mathcal{Z}u$ by

$$\mathcal{Z}u := Z \cdot \nabla u - ik\beta u + \alpha u,$$

where Z , α , β are real-valued, with $Z \cdot n \geq 0$ on the boundary, and

$$2\Re \int_{\Omega_R} \overline{\mathcal{Z}u} f \, dx = 2\Re \int_{\Omega_R} \overline{\mathcal{Z}u} (\Delta u + k^2 u) \, dx = \int_{\Omega_R} +ve + \int_{\partial\Omega_R} +ve$$

In **rough surface scattering** (C-W, Monk 2005) use $Z(x) = x_d e_d$, $\alpha = 1/2$, $\beta = R$, to get

$$\int_{\Omega_R} |\partial_d u|^2 \, dx \leq -2\Re \int_{\Omega_R} \overline{\mathcal{Z}u} f \, dx \leq \epsilon \|\mathcal{Z}u\|_{L^2(\Omega_R)}^2 + \epsilon^{-1} \|f\|_{L^2(\Omega_R)}^2;$$

then use Friedrichs inequality to bound $\|u\|_{L^2(\Omega_R)}$ in terms of $\|\partial_d u\|_{L^2(\Omega_R)}$.

How is our new estimate for parabolic trapping proved?

Define **Morawetz multiplier** $\mathcal{Z}u$ by

$$\mathcal{Z}u := Z \cdot \nabla u - ik\beta u + \alpha u,$$

where Z , α , β are real-valued, with $Z \cdot n \geq 0$ on the boundary, and

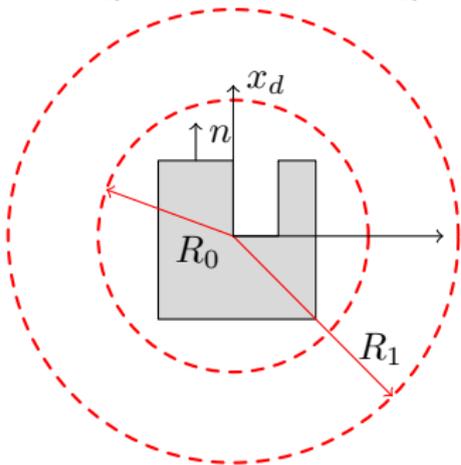
$$2\Re \int_{\Omega_R} \overline{\mathcal{Z}u} f \, dx = 2\Re \int_{\Omega_R} \overline{\mathcal{Z}u} (\Delta u + k^2 u) \, dx = \int_{\Omega_R} +ve + \int_{\partial\Omega_R} +ve + \int_{\Omega_R} \text{small}$$

Define **Morawetz multiplier** $\mathcal{Z}u$ by

$$\mathcal{Z}u := Z \cdot \nabla u - ik\beta u + \alpha u,$$

where Z , α , β are real-valued, with $Z \cdot n \geq 0$ on the boundary, and

$$2\Re \int_{\Omega_R} \overline{\mathcal{Z}u} f \, dx = 2\Re \int_{\Omega_R} \overline{\mathcal{Z}u} (\Delta u + k^2 u) \, dx = \int_{\Omega_R} +ve + \int_{\partial\Omega_R} +ve + \int_{\Omega_R} \text{small}$$



Choose $R_1 > R_0 > 0$ and set

$$Z(x) = x_d e_d \text{ for } |x| \leq R_0, \quad Z(x) = x \text{ for } |x| \geq R_1.$$

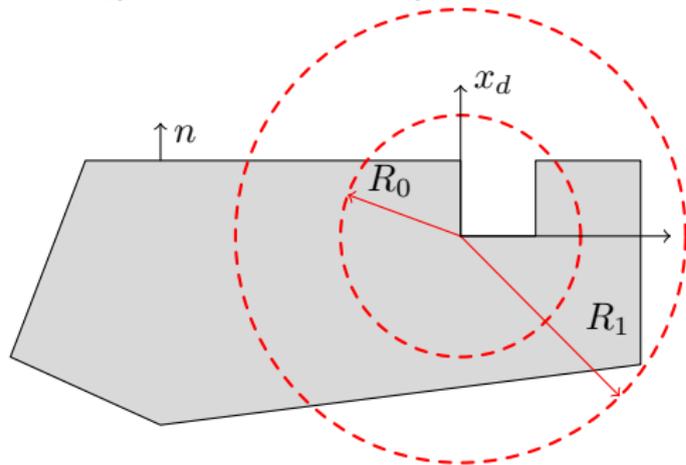
Resolvent estimate obtained if $Z \cdot n = x_d e_d \cdot n \geq 0$ on boundary & $R_1/R_0 \geq 121$.

Define **Morawetz multiplier** $\mathcal{Z}u$ by

$$\mathcal{Z}u := Z \cdot \nabla u - ik\beta u + \alpha u,$$

where Z , α , β are real-valued, with $Z \cdot n \geq 0$ on the boundary, and

$$2\Re \int_{\Omega_R} \overline{\mathcal{Z}u} f \, dx = 2\Re \int_{\Omega_R} \overline{\mathcal{Z}u} (\Delta u + k^2 u) \, dx = \int_{\Omega_R} +ve + \int_{\partial\Omega_R} +ve + \int_{\Omega_R} \text{small}$$



Choose $R_1 > R_0 > 0$ and set

$$Z(x) = x_d e_d \text{ for } |x| \leq R_0, \quad Z(x) = x \text{ for } |x| \geq R_1.$$

Resolvent estimate obtained if $Z \cdot n \geq 0$ on boundary and $R_1/R_0 \geq 121$.

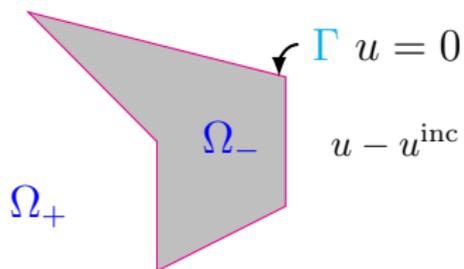
Overview of Talk

- 1 Helmholtz equation: what is it and why interesting
- 2 What is this talk about?
- 3 Resolvent estimates
 - What are they?
 - The three known estimates and their geometries
 - A new estimate for parabolic trapping
 - The Morawetz/Rellich identity method of proof
- 4 Implications for Boundary Integral Equations
- 5 Implications for hp -BEM
- 6 Conclusions

Integral Equations and k -Explicit Bounds

$\mathcal{W} \rightarrow u^{\text{inc}}$

$$\Delta u + k^2 u = 0$$



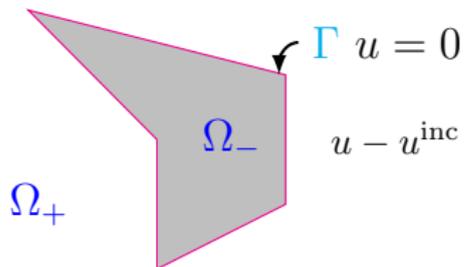
$u - u^{\text{inc}}$ satisfies radiation condition

Assume throughout that Ω_- is bounded and Lipschitz.

Integral Equations and k -Explicit Bounds

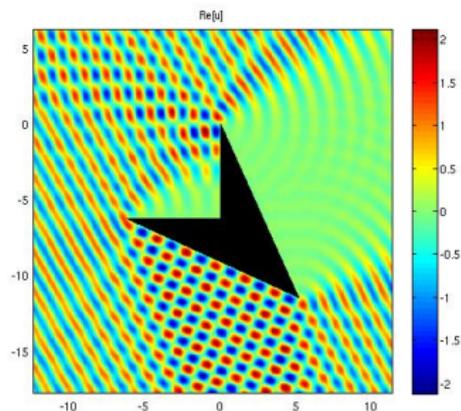
$\mathcal{W} \rightarrow u^{\text{inc}}$

$$\Delta u + k^2 u = 0$$

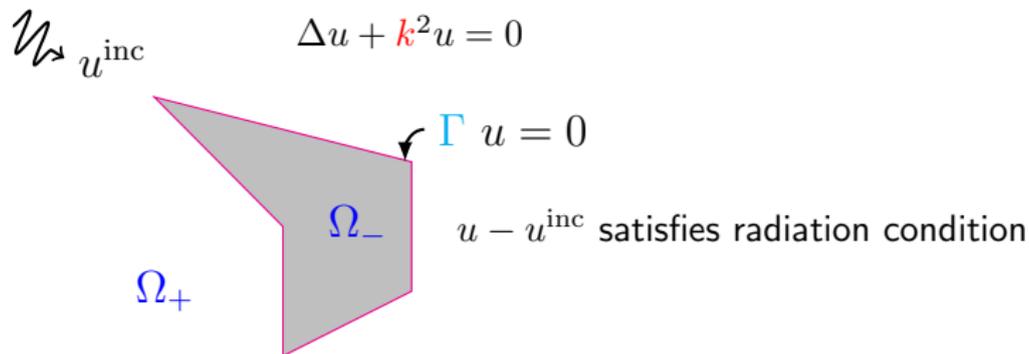


$u - u^{\text{inc}}$ satisfies radiation condition

Assume throughout that Ω_- is bounded and Lipschitz. Plot of $\Re(u(x)) = U(x, 0)$:



Integral Equations and k -Explicit Bounds



Assume throughout that Ω_- is bounded and Lipschitz.

Theorem (Green's Representation Theorem)

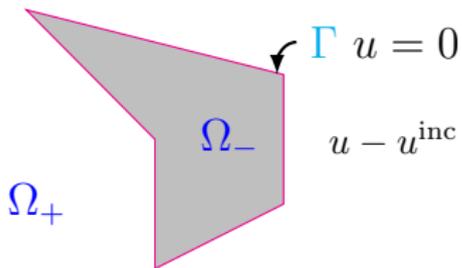
$$u(x) = u^{\text{inc}}(x) + \int_{\Gamma} \Phi(x, y) \partial_n^+ u(y) ds(y), \quad x \in \Omega_+.$$

where

$$\Phi(x, y) := \frac{i}{4} H_0^{(1)}(k|x-y|) \quad (2\text{D}), \quad := \frac{1}{4\pi} \frac{e^{ik|x-y|}}{|x-y|} \quad (3\text{D}).$$

$\mathcal{W} \rightarrow u^{\text{inc}}$

$$\Delta u + k^2 u = 0$$



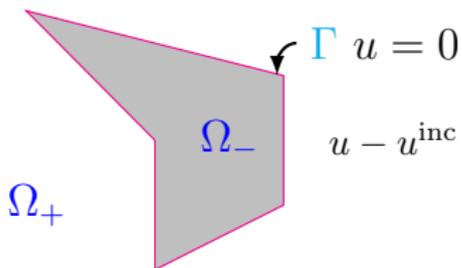
$u - u^{\text{inc}}$ satisfies radiation condition

Theorem (Green's Representation Theorem)

$$u(x) = u^{\text{inc}}(x) + \int_{\Gamma} \Phi(x, y) \partial_n^+ u(y) ds(y), \quad x \in \Omega_+.$$

$\mathcal{W}_\rightarrow u^{\text{inc}}$

$$\Delta u + k^2 u = 0$$



$$\Gamma \quad u = 0$$

$u - u^{\text{inc}}$ satisfies radiation condition

 Ω_+ Ω_-

Theorem (Green's Representation Theorem)

$$u(x) = u^{\text{inc}}(x) + \int_{\Gamma} \Phi(x, y) \partial_n^+ u(y) ds(y), \quad x \in \Omega_+.$$

Taking a linear combination of Dirichlet (γ_+) and Neumann (∂_n^+) traces, we get the **boundary integral equation** (Burton & Miller 1971)

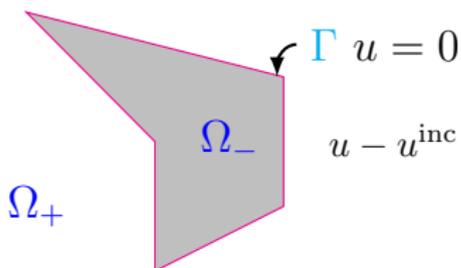
$$\frac{1}{2} \partial_n^+ u(x) + \int_{\Gamma} \left(\frac{\partial \Phi(x, y)}{\partial n(x)} + i\eta \Phi(x, y) \right) \partial_n^+ u(y) ds(y) = f(x), \quad x \in \Gamma,$$

where

$$f := \partial_n^+ u^{\text{inc}} + i\eta \gamma_+ u^{\text{inc}}.$$

$\mathcal{W} \rightarrow u^{\text{inc}}$

$$\Delta u + k^2 u = 0$$



$u - u^{\text{inc}}$ satisfies radiation condition

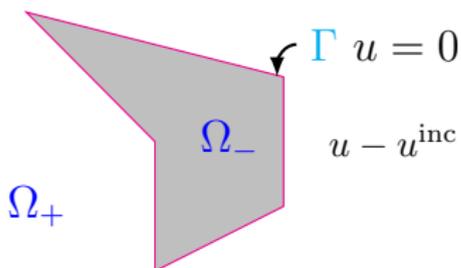
$$\frac{1}{2} \partial_n^+ u(x) + \int_{\Gamma} \left(\frac{\partial \Phi(x, y)}{\partial n(x)} + i\eta \Phi(x, y) \right) \partial_n^+ u(y) ds(y) = f(x), \quad x \in \Gamma,$$

in operator form

$$A_{k, \eta} \partial_n^+ u = f := \partial_n^+ u^{\text{inc}} + i\eta \gamma_+ u^{\text{inc}}.$$

$\mathcal{W} \rightarrow u^{\text{inc}}$

$$\Delta u + k^2 u = 0$$



$u - u^{\text{inc}}$ satisfies radiation condition

$$\frac{1}{2} \partial_n^+ u(x) + \int_{\Gamma} \left(\frac{\partial \Phi(x, y)}{\partial n(x)} + i\eta \Phi(x, y) \right) \partial_n^+ u(y) ds(y) = f(x), \quad x \in \Gamma,$$

in operator form

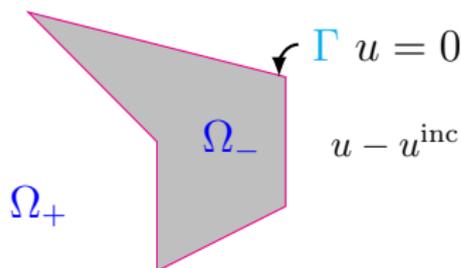
$$A_{k, \eta} \partial_n^+ u = f := \partial_n^+ u^{\text{inc}} + i\eta \gamma_+ u^{\text{inc}}.$$

Theorem (Burton & Miller 1971, Mitrea 1996, C-W & Langdon 2007)

If $\eta \in \mathbb{R}$, $\eta \neq 0$, then this integral equation is uniquely solvable in $L^2(\Gamma)$.

$\mathcal{W} \rightarrow u^{\text{inc}}$

$$\Delta u + k^2 u = 0$$

 $u - u^{\text{inc}}$ satisfies radiation condition

$$\frac{1}{2} \partial_n^+ u(x) + \int_{\Gamma} \left(\frac{\partial \Phi(x, y)}{\partial n(x)} + i\eta \Phi(x, y) \right) \partial_n^+ u(y) ds(y) = f(x), \quad x \in \Gamma,$$

in operator form

$$A_{k, \eta} \partial_n^+ u = f := \partial_n^+ u^{\text{inc}} + i\eta \gamma_+ u^{\text{inc}}.$$

Theorem (Burton & Miller 1971, Mitrea 1996, C-W & Langdon 2007)

If $\eta \in \mathbb{R}$, $\eta \neq 0$, then this integral equation is uniquely solvable in $L^2(\Gamma)$.

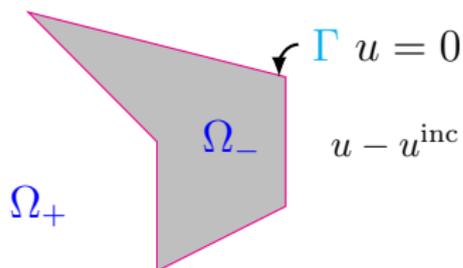
The standard choice is $\eta = k$, and with this choice we have

$$\|A_{k, k}^{-1}\|_{L^2(\Gamma) \rightarrow L^2(\Gamma)} \lesssim 1$$

if Ω_- is **star-shaped** (C-W, Monk 2008) or C^∞ and **nontrapping** (Baskin, Spence, Wunsch 2016).

$\mathcal{W} \rightarrow u^{\text{inc}}$

$$\Delta u + k^2 u = 0$$

 $u - u^{\text{inc}}$ satisfies radiation condition

$$\frac{1}{2} \partial_n^+ u(x) + \int_{\Gamma} \left(\frac{\partial \Phi(x, y)}{\partial n(x)} + i\eta \Phi(x, y) \right) \partial_n^+ u(y) ds(y) = f(x), \quad x \in \Gamma,$$

in operator form

$$A_{k, \eta} \partial_n^+ u = f := \partial_n^+ u^{\text{inc}} + i\eta \gamma_+ u^{\text{inc}}.$$

Theorem (Burton & Miller 1971, Mitrea 1996, C-W & Langdon 2007)

If $\eta \in \mathbb{R}$, $\eta \neq 0$, then this integral equation is uniquely solvable in $L^2(\Gamma)$.

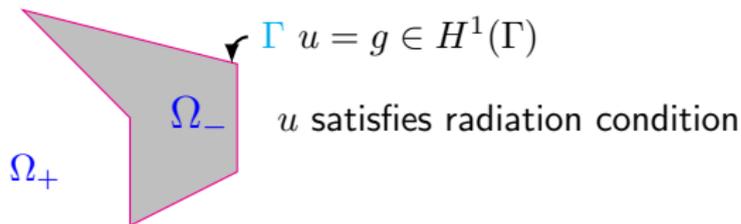
The standard choice is $\eta = k$, and with this choice we have

$$\|A_{k, k}^{-1}\|_{L^2(\Gamma) \rightarrow L^2(\Gamma)} \lesssim 1$$

if Ω_- is **star-shaped** (C-W, Monk 2008) or C^∞ and **nontrapping** (Baskin, Spence, Wunsch 2016). **But what if Ω_- is trapping?**

A recipe for bounding $\|A_{k,k}^{-1}\|$ (C-W, Spence, Gibbs, Smyshlyaev 2017)

$$\Delta u + k^2 u = f \in L^2(\Omega_+), \text{ compactly supported}$$

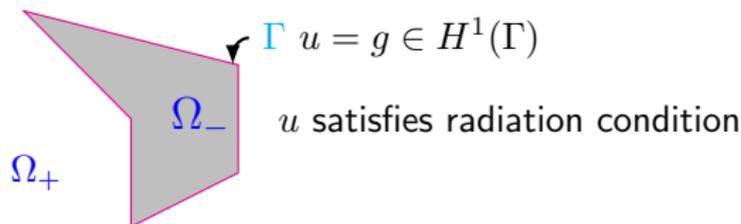


$$\Gamma u = g \in H^1(\Gamma)$$

u satisfies radiation condition

A recipe for bounding $\|A_{k,k}^{-1}\|$ (C-W, Spence, Gibbs, Smyshlyaev 2017)

$$\Delta u + k^2 u = f \in L^2(\Omega_+), \text{ compactly supported}$$



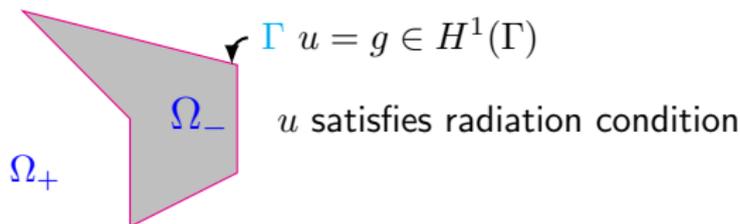
Step 1 (Resolvent Estimate). Show that, for every $R > 0$, if $g = 0$,

$$\|\nabla u\|_{L^2(\Omega_R)} + k\|u\|_{L^2(\Omega_R)} \lesssim c(k)\|f\|_{L^2(\Omega_+)},$$

where $\Omega_R := \{x \in \Omega_+ : |x| < R\}$.

A recipe for bounding $\|A_{k,k}^{-1}\|$ (C-W, Spence, Gibbs, Smyshlyaev 2017)

$$\Delta u + k^2 u = f \in L^2(\Omega_+), \text{ compactly supported}$$



Step 1 (Resolvent Estimate). Show that, for every $R > 0$, if $g = 0$,

$$\|\nabla u\|_{L^2(\Omega_R)} + k\|u\|_{L^2(\Omega_R)} \lesssim c(k)\|f\|_{L^2(\Omega_+)},$$

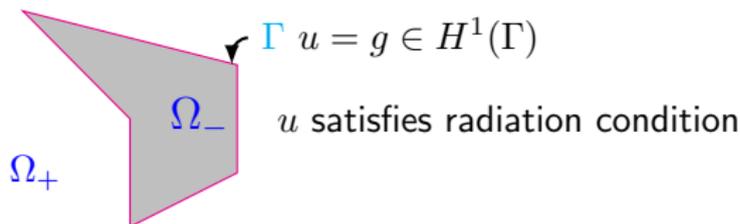
where $\Omega_R := \{x \in \Omega_+ : |x| < R\}$.

Step 2 (DtN Map Bound). It follows that, if $f = 0$,

$$\|\partial_n^+ u\|_{L^2(\Gamma)} \lesssim c(k) (\|\nabla_\Gamma g\|_{L^2(\Gamma)} + k\|g\|_{L^2(\Gamma)})$$

A recipe for bounding $\|A_{k,k}^{-1}\|$ (C-W, Spence, Gibbs, Smyshlyaev 2017)

$$\Delta u + k^2 u = f \in L^2(\Omega_+), \text{ compactly supported}$$



Step 1 (Resolvent Estimate). Show that, for every $R > 0$, if $g = 0$,

$$\|\nabla u\|_{L^2(\Omega_R)} + k\|u\|_{L^2(\Omega_R)} \lesssim c(k)\|f\|_{L^2(\Omega_+)},$$

where $\Omega_R := \{x \in \Omega_+ : |x| < R\}$.

Step 2 (DtN Map Bound). It follows that, if $f = 0$,

$$\|\partial_n^+ u\|_{L^2(\Gamma)} \lesssim c(k) (\|\nabla_{\Gamma} g\|_{L^2(\Gamma)} + k\|g\|_{L^2(\Gamma)})$$

Step 3 As (C-W, Graham, Langdon, Spence 2012)

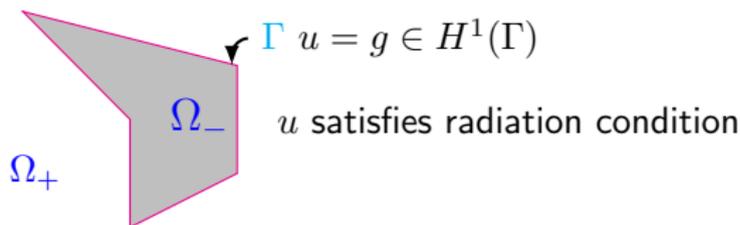
$$A_{k,k}^{-1} = I - (P_{DtN}^+ - ik)P_{ItD}^-$$

and P_{ItD}^- is bounded in Spence (2015), Baskin, Spence, Wunsch (2016), it follows that

$$\|A_{k,k}^{-1}\|_{L^2(\Gamma) \rightarrow L^2(\Gamma)} \lesssim c(k)k^{1/2}$$

A recipe for bounding $\|A_{k,k}^{-1}\|$ (C-W, Spence, Gibbs, Smyshlyaev 2017)

$$\Delta u + k^2 u = f \in L^2(\Omega_+), \text{ compactly supported}$$



Step 1 (Resolvent Estimate). Show that, for every $R > 0$, if $g = 0$,

$$\|\nabla u\|_{L^2(\Omega_R)} + k\|u\|_{L^2(\Omega_R)} \lesssim c(k)\|f\|_{L^2(\Omega_+)},$$

where $\Omega_R := \{x \in \Omega_+ : |x| < R\}$.

Step 2 (DtN Map Bound). It follows that, if $f = 0$,

$$\|\partial_n^+ u\|_{L^2(\Gamma)} \lesssim c(k) (\|\nabla_{\Gamma} g\|_{L^2(\Gamma)} + k\|g\|_{L^2(\Gamma)})$$

Step 3 As (C-W, Graham, Langdon, Spence 2012)

$$A_{k,k}^{-1} = I - (P_{DtN}^+ - ik)P_{ItD}^-$$

and P_{ItD}^- is bounded in Spence (2015), Baskin, Spence, Wunsch (2016), it follows that

$$\|A_{k,k}^{-1}\|_{L^2(\Gamma) \rightarrow L^2(\Gamma)} \lesssim c(k)$$

if each component of Ω_- is star-shaped or C^∞ .

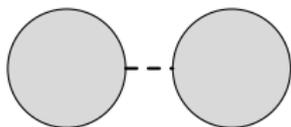
Recap of resolvent estimates for trapping obstacles

$$\|\nabla u\|_{L^2(\Omega_R)} + k\|u\|_{L^2(\Omega_R)} \lesssim c(k)\|f\|_{L^2(\Omega_R)}, \quad \text{for } k \geq k_0 > 0,$$

where $c(k) = 1$ for **nontrapping** obstacles, and



$c(k) = \exp(\alpha k)$
elliptic



$c(k) = \log(2 + k)$
hyperbolic

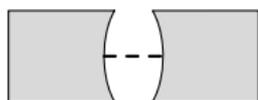


$c(k) = k^2$
parabolic

Recap of resolvent estimates for trapping obstacles

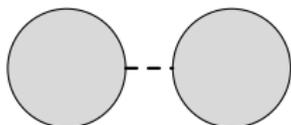
$$\|\nabla u\|_{L^2(\Omega_R)} + k\|u\|_{L^2(\Omega_R)} \lesssim c(k)\|f\|_{L^2(\Omega_R)}, \quad \text{for } k \geq k_0 > 0,$$

where $c(k) = 1$ for **nontrapping** obstacles, and



$$c(k) = \exp(\alpha k)$$

elliptic



$$c(k) = \log(2 + k)$$

hyperbolic



$$c(k) = k^2$$

parabolic

Applying our general recipe

$$\|A_{k,k}^{-1}\|_{L^2(\Gamma) \rightarrow L^2(\Gamma)} \lesssim c(k)$$

Recap of resolvent estimates for trapping obstacles

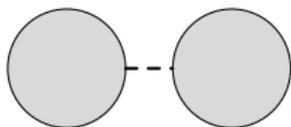
$$\|\nabla u\|_{L^2(\Omega_R)} + k\|u\|_{L^2(\Omega_R)} \lesssim c(k)\|f\|_{L^2(\Omega_R)}, \quad \text{for } k \geq k_0 > 0,$$

where $c(k) = 1$ for **nontrapping** obstacles, and



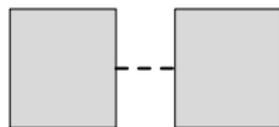
$$c(k) = \exp(\alpha k)$$

elliptic



$$c(k) = \log(2 + k)$$

hyperbolic



$$c(k) = k^2$$

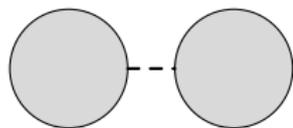
parabolic

Applying our general recipe, for some $N \geq 0$,

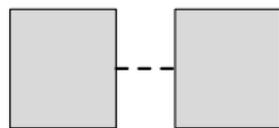
$$\|A_{k,k}^{-1}\|_{L^2(\Gamma) \rightarrow L^2(\Gamma)} \lesssim c(k) \lesssim k^N$$

in the **nontrapping** and **hyperbolic** and **parabolic trapping** cases.

Application to hp -BEM analysis



hyperbolic



parabolic

For these configurations $\exists N \geq 0$ s.t. $\|A_{k,k}^{-1}\|_{L^2(\Gamma) \rightarrow L^2(\Gamma)} \lesssim k^N$, $k \geq k_0 > 0$.

Corollary (Löhndorf, Melenk 2011)

Suppose Γ is analytic and \mathcal{T}_h is a quasi-uniform triangulation with mesh size h . Then, given $k_0 > 0$, $\exists C_1, C_2, C_3$ such that, if $k \geq k_0$,

$$\frac{kh}{p} \leq C_1, \quad \text{and} \quad p \geq C_2 \log(2 + k),$$

then the Galerkin hp -BEM solution $v_{hp} \in \mathcal{S}^p(\mathcal{T}_h)$ satisfies the quasi-optimal error estimate

$$\|v_{hp} - \partial_n^+ u\|_{L^2(\Gamma)} \leq C_3 \inf_{v \in \mathcal{S}^p(\mathcal{T}_h)} \|v - \partial_n^+ u\|_{L^2(\Gamma)}.$$

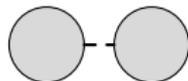
Conclusions

In this talk you have seen:

- All the resolvent estimates that exist for (Dirichlet) obstacles



elliptic



hyperbolic



parabolic

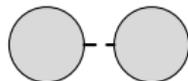
Conclusions

In this talk you have seen:

- All the resolvent estimates that exist for (Dirichlet) obstacles



elliptic



hyperbolic



parabolic

- In particular our new bound for parabolic trapping obstacles

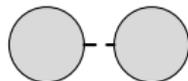
Conclusions

In this talk you have seen:

- All the resolvent estimates that exist for (Dirichlet) obstacles



elliptic



hyperbolic



parabolic

- In particular our new bound for parabolic trapping obstacles
- The Morawetz/Rellich identity method for proving these estimates

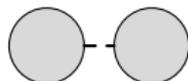
Conclusions

In this talk you have seen:

- All the resolvent estimates that exist for (Dirichlet) obstacles



elliptic



hyperbolic



parabolic

- In particular our new bound for parabolic trapping obstacles
- The Morawetz/Rellich identity method for proving these estimates
- How resolvent estimates lead in a “black box” way to:
 - bounds on the DtN map
 - bounds on $\|A_{k,k}^{-1}\|_{L^2 \rightarrow L^2}$
 - hp -BEM quasi-optimality

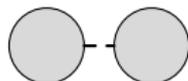
Conclusions

In this talk you have seen:

- All the resolvent estimates that exist for (Dirichlet) obstacles



elliptic



hyperbolic



parabolic

- In particular our new bound for parabolic trapping obstacles
- The Morawetz/Rellich identity method for proving these estimates
- How resolvent estimates lead in a “black box” way to:
 - bounds on the DtN map
 - bounds on $\|A_{k,k}^{-1}\|_{L^2 \rightarrow L^2}$
 - hp -BEM quasi-optimality

Not covered today are k -explicit results for h -BEM, FEM, and bounds on $A_{k,k}^{-1}$ as an operator on $H^s(\Gamma)$, for $-1 \leq s \leq 0$.

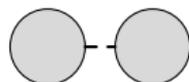
Conclusions

In this talk you have seen:

- All the resolvent estimates that exist for (Dirichlet) obstacles



elliptic



hyperbolic



parabolic

- In particular our new bound for parabolic trapping obstacles
- The Morawetz/Rellich identity method for proving these estimates
- How resolvent estimates lead in a “black box” way to:
 - bounds on the DtN map
 - bounds on $\|A_{k,k}^{-1}\|_{L^2 \rightarrow L^2}$
 - hp -BEM quasi-optimality

Not covered today are k -explicit results for h -BEM, FEM, and bounds on $A_{k,k}^{-1}$ as an operator on $H^s(\Gamma)$, for $-1 \leq s \leq 0$.

More details see:

C-W, Spence, Gibbs, Smyshlyaev 2017, *High-frequency bounds for the Helmholtz equation under parabolic trapping and applications in numerical analysis*, arXiv:1708.08415